Paulo Santos presents...

# Astrophysical constraints on Ungravity inspired models

In collaboration with Orfeu Bertolami and Jorge Páramos

Institute of Theoretical Astrophysics,
University of Oslo

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# Astrophysical constraints on Ungravity inspired models

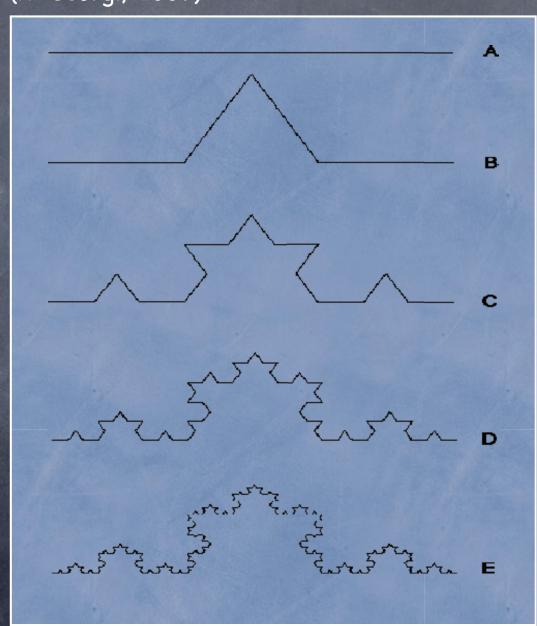
Based on work with the same title presented in Physical Review D 80, 022001 (2009)

[ arXiv:0905.1602v2 [astro-ph.SR] ]

### What are Unparticles?

(H. Georgi, 2007)

Massive ScaleInvariantObjects



### How to interpret Unparticles?

Non-integer number of particles

Tower of continuous mass fields

### Ungravity

(H. Goldberg and P. Nath, 2007)

$$V(r) = -\frac{G_U M m}{r} \left[ 1 + \left(\frac{R_G}{r}\right)^{2d_U - 2} \right]$$

with:

$$R_G = \frac{1}{\pi \Lambda_U} \left(\frac{M_{Pl}}{M_*}\right)^{1/(d_U - 1)} \times \left[\frac{2(2 - \alpha)}{\pi} \frac{\Gamma(d_U + \frac{1}{2})\Gamma(d_U - \frac{1}{2})}{\Gamma(2d_U)}\right]^{1/(2d_U - 2)}$$

# Ungravity matching the Newtonian potential

(O. Bertolami, J. Páramos and P. Santos, 2009)

To match the Newtonian potential for the non-interacting cases –  $d_U = 1$ ,  $R_G = 0$  (for  $d_U > 1$ ) and  $R_G = \infty$  (for  $d_U < 1$ ) – we have:

$$G_U = \frac{G}{1 + \left(\frac{R_G}{R_0}\right)^{2d_U - 2}}$$

For du close to unity:

$$V(r) = -\frac{GMm}{2r} \left[ 1 + \left(\frac{R_G}{r}\right)^{2d_U - 2} \right]$$

### How to test Ungravity?

(O. Bertolami, J. Páramos and P. Santos, 2009)

Astrophysical constraints using a modified Lane-Emden equation applied to the Sun.

### Usual Lane-Emden Equation

(Wikipedia, for example...)

Hydrostatic Equilibrium:

$$\frac{d}{dr}\left(\frac{r^2}{\rho}\frac{dP}{dr}\right) = -G\frac{dM(r)}{dr}$$

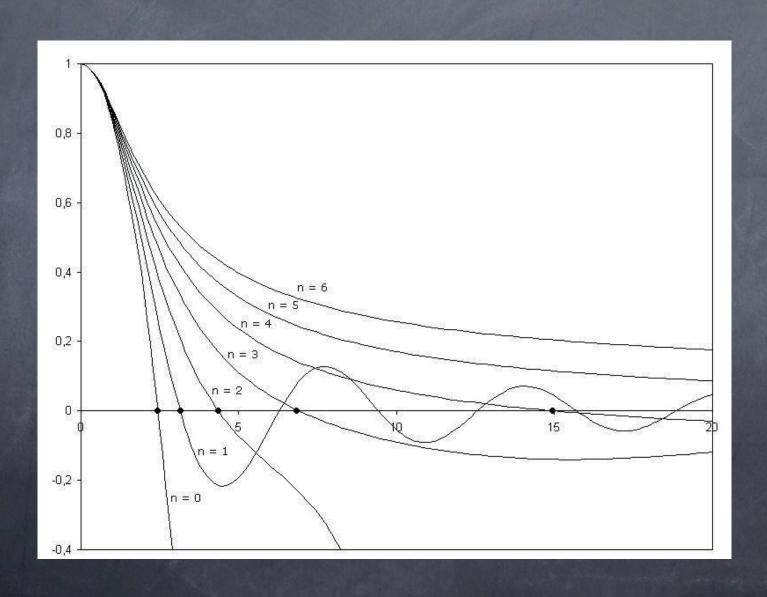
Mass Conservation:

$$\frac{dM(r)}{dr} = 4\pi\rho(r)r^2$$

Lane-Emden Equation:

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta_0}{d\xi} \right) = -\theta_0^n$$

### A plot of $\Theta$ as a function of $\xi$



# Modified Lane-Emden Equation (I)

(O. Bertolami, J. Páramos and P. Santos, 2009)

Hydrostatic Equilibrium:

$$\frac{r^2}{\rho} \frac{dP(r)}{dr} = -\frac{GM(r)}{2} \left[ 1 + (2d_U - 1) \left( \frac{R_G}{r} \right)^{2d_U - 2} \right]$$

Mass Conservation:

$$M(\xi) = -4\pi \left(\frac{(n+1)K}{2\pi G}\right)^{3/2} \rho_c^{(3-n)/2n} \xi^2 \frac{d\theta}{d\xi}$$

# Modified Lane-Emden Equation (II)

(O. Bertolami, J. Páramos and P. Santos, 2009)

#### Modified Lane-Emden Equation:

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{\theta^n}{\xi^2} \left( \frac{\xi_G}{\xi} \right) = -\frac{\theta^n}{2} \left[ 1 + (2d_U - 1) \left( \frac{\xi_G}{\xi} \right)^{2d_U - 2} \right] - \frac{\theta^n}{\xi^2} \left( \frac{\xi_G}{\xi} \right) = -\frac{\theta^n}{2} \left[ \frac{\xi_G}{\xi} \right] + \frac{\theta^n}{\xi^2} \left( \frac{\xi_G}{\xi} \right) = -\frac{\theta^n}{2} \left[ \frac{\xi_G}{\xi} \right] + \frac{\theta^n}{\xi^2} \left( \frac{\xi_G}{\xi} \right) + \frac{\theta^n}{\xi^2} \left( \frac{$$

$$(2d_U - 1)(d_U - 1)\frac{1}{\xi}\frac{d\theta}{d\xi}\left(\frac{\xi_G}{\xi}\right)^{2d_U - 2}$$

### Temperature Method

(O. Bertolami, J. Páramos and P. Santos, 2009)

 $\circ$  6% is the observational uncertainty for the Sun's temperature -  $|T_r-1|=6\%$ 

Newtonian Temperature / Ungravity Temperature:

$$T_{c0} \propto \left[ \xi_{10} \left( \frac{d\theta_0}{d\xi} \right)_{\xi=\xi_{10}} \right]^{-1} , \quad T_c \propto \left[ \xi_1 \left( \frac{d\theta}{d\xi} \right)_{\xi=\xi_1} \right]^{-1}$$

Ratio:

$$T_r \equiv \frac{T_c}{T_{c0}} = \frac{\xi_{10}}{\xi_1} \frac{\frac{d\theta_0}{d\xi}}{\frac{d\theta}{d\xi}}$$

### Lower Bound on M\*

(O. Bertolami, J. Páramos and P. Santos, 2009)

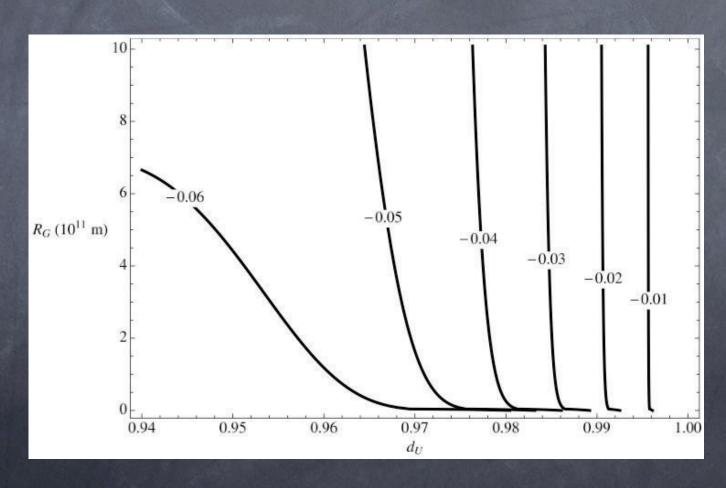
$$\frac{M_*}{M_{Pl}} > \left[\pi \Lambda_U R_6(d_U)\right]^{1-d_U} f(d_U, \alpha)$$

with:

$$f(d_U, \alpha) = \sqrt{\frac{2(2-\alpha)}{\pi} \frac{\Gamma(d_U + \frac{1}{2})\Gamma(d_U - \frac{1}{2})}{\Gamma(2d_U)}}$$

# Results Constraining $d_U$ and $R_G$ (I)

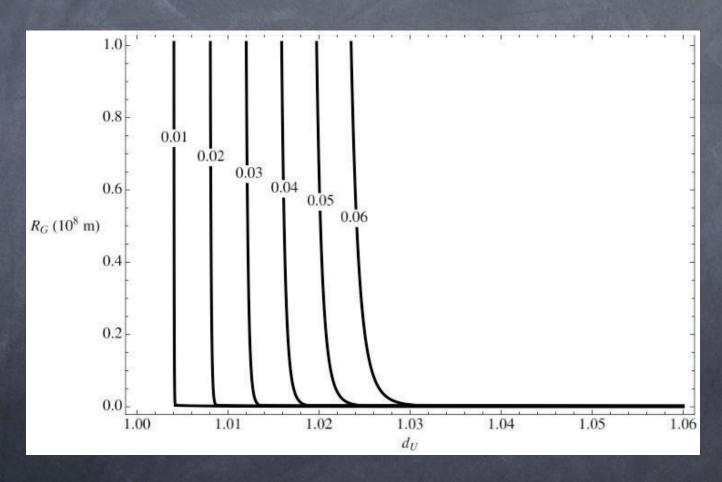
(O. Bertolami, J. Páramos and P. Santos, 2009)



Contour plot of  $T_{r}-1$  as a function of log  $R_{G}$  and  $d_{U}$ 

# Results Constraining $d_U$ and $R_G$ (II)

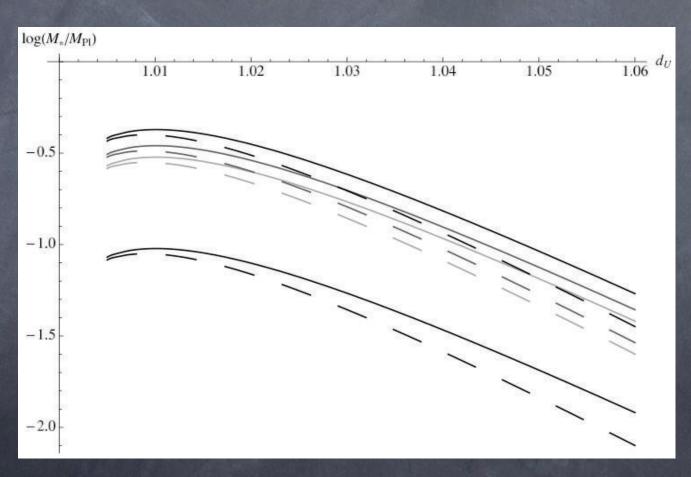
(O. Bertolami, J. Páramos and P. Santos, 2009)



Contour plot of  $T_{r}-1$  as a function of log  $R_{G}$  and  $d_{U}$ 

# Results Constraining M\* (I)

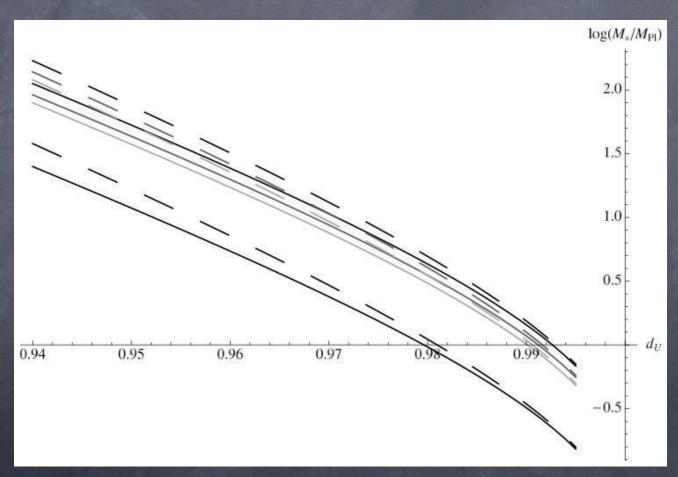
(O. Bertolami, J. Páramos and P. Santos, 2009)



Lower bound of log(M\*/M<sub>Pl</sub>) for  $\alpha$ =0 (black),  $\alpha$ =2/3 (dark grey),  $\alpha$ =1 (light grey),  $\alpha$ =1.9 (black, lower curve) and  $\Lambda_U$ =1 TeV (solid),  $\Lambda_U$ =10<sup>3</sup> TeV (dashed)

# Results Constraining M\* (II)

(O. Bertolami, J. Páramos and P. Santos, 2009)



Lower bound of log(M\*/MPI) for  $\alpha$ =0 (black),  $\alpha$ =2/3 (dark grey),  $\alpha$ =1 (light grey),  $\alpha$ =1.9 (black, lower curve) and  $\Lambda_U$ =1 TeV (solid),  $\Lambda_U$ =10<sup>3</sup> TeV (dashed)

### Previous Constraints

Method	Range of $d_U$	Lower bound on $M_*$
Astrophysical Constraints [1]	1 - 2	$10^{19} \text{ TeV } (1 \leqslant \Lambda_U \leqslant 10^3 \text{ TeV})$
Cosmological Constraints [2]	1.1- 2	$10 - 10^3  \mathrm{TeV}  \left( \Lambda_U = 1  \mathrm{TeV} \right)$
Ungravity Constraints [3]	2 - 2.2	$10^3 - 10^{10} \text{ TeV } (1 \leqslant \Lambda_U \leqslant 10^3 \text{ TeV})$

[1] S. Das, S. Mohanty and K. Rao, 2008[2] J. McDonald, 2009[3] H. Goldberg and P. Nath, 2009

# Conclusions (I)

Our method allows for obtaining lower bounds for M\*:  $(10^{-2} - 10^{-1}) \text{ M}_{Pl} \text{ for } d_{U} ≈ 1 \text{ and } (10^{-1} - 10^{2}) \text{ M}_{Pl} \text{ for } d_{U} ≈ 1 \text{, in both cases for } Λ_{U} = 1 \text{ TeV, } 10^{3} \text{ TeV}$ 

The result for  $d_{U} \gtrsim 1$  is at least as stringent as those previously obtained

# Conclusions (II)

The d<sub>U</sub> ≤ 1 case was not examined so far

 $\ \ \, \ \ \,$  The objection for  $d_U \lesssim 1$  of quick divergence is only true for  $d_U \ll 1$ 

$$(r/R_G)^{2-2d_U} \approx 1 + 2(d_U - 1)\log(r/R_g)$$