

Models of Dark Energy Accretion onto Black Holes

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- Black hole interaction with dark energy
[Babichev et al., 2005, Guariento et al., 2008]



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- Dark energy models and evolution



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- Black hole interaction with dark energy
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- Dark energy models and evolution
- Black hole and dark energy thermodynamics (negativity of entropy \times negativity of temperature \times neither) [Lima et al., 2010]



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- Black hole and dark energy thermodynamics (negativity of entropy \times negativity of temperature \times neither) [Lima et al., 2010]
- Contribution of black holes to the evolution of the universe
[Khlopov, 2007]



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- Contribution of black holes to the evolution of the universe
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- Back reaction analysis [Faraoni and Jacques, 2007]



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- 4-momentum in a box of volume V



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- 4-momentum in a box of volume V

$$p^\mu = \int_V T^{\mu\nu} d\Sigma_\nu$$



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$$p^\mu = \int_V T^{\mu\nu} d\Sigma_\nu = VT^{\mu\nu} u_\nu$$



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- 4-momentum transferred from the box surface S during $\Delta\tau$

$$\Delta p^\mu = S\Delta\tau T^{\mu\nu} \sigma_\nu$$



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- Energy variation through the horizon of a Schwarzschild black hole

$$\frac{dE_{\text{inside}}}{d\tau} = \frac{dm}{d\tau}$$



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$$\frac{dE_{\text{inside}}}{d\tau} = \frac{dm}{d\tau} = \frac{dE_{\text{outside}}}{d\tau} = ST^{\mu\nu} u_\mu \sigma_\nu$$



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$$\frac{dE_{\text{inside}}}{d\tau} = \frac{dm}{d\tau} = \frac{dE_{\text{outside}}}{d\tau} = ST^{\mu\nu} u_\mu \sigma_\nu$$

$$\boxed{\frac{dm}{dt} = ST_0^1} \quad (1)$$



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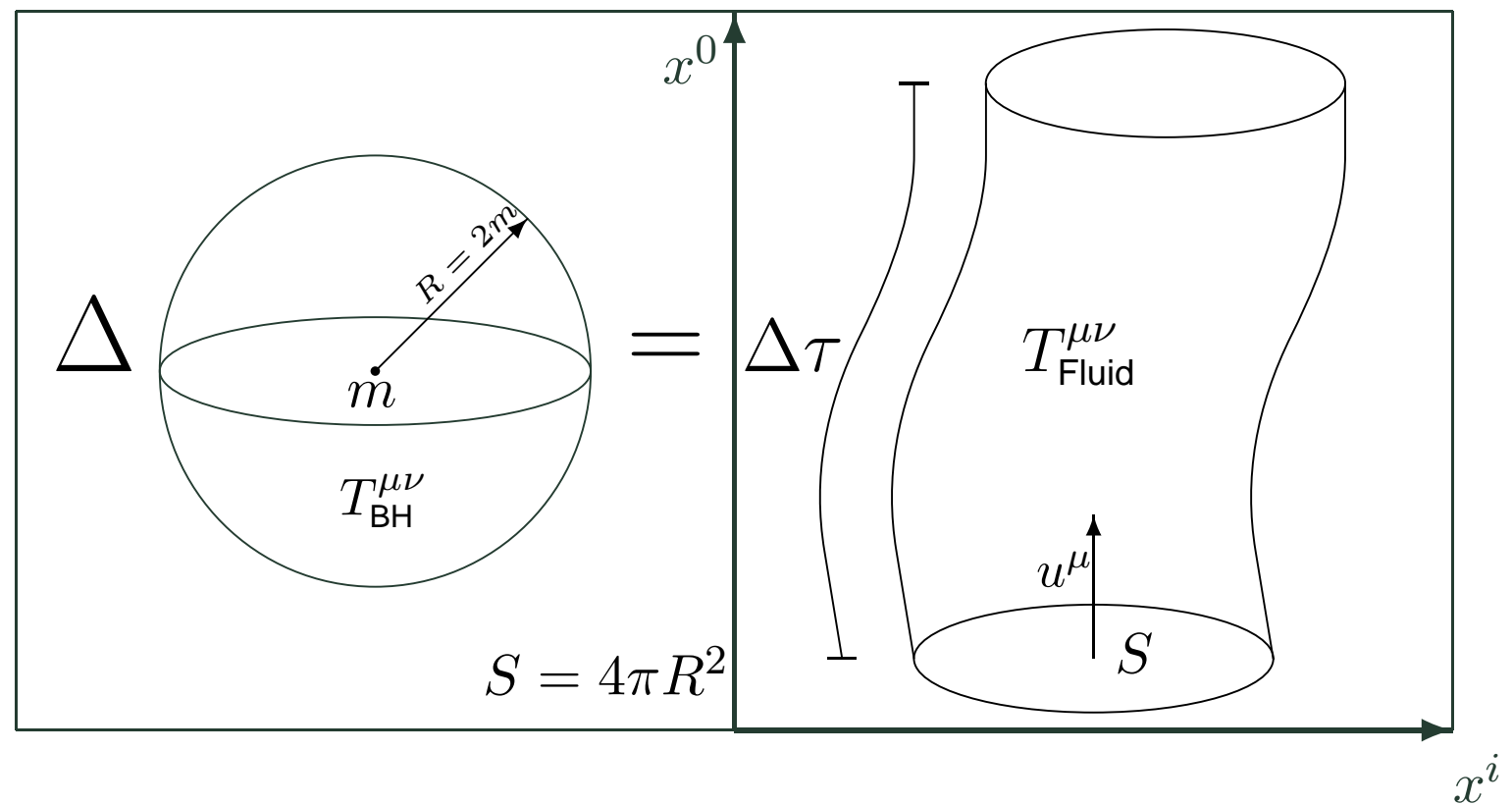
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$$T^{\mu\nu}_{;\nu} = 0;$$



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$$T^{\mu\nu}{}_{;\nu} = 0; \quad u_{\mu}T^{\mu\nu}{}_{;\nu} = 0 \quad (2)$$



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- Accretion of dark energy and radiation with Hawking evaporation

$$\frac{dm}{dt} = -\frac{A(m)}{m^2} + m^2 [27\pi\rho_{\text{rad}}(T) + 16\pi(1+w)\rho_{\text{DE}}]$$



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- Phantom dark energy causes regime transitions



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- Phantom dark energy causes regime transitions

- Radiation accretion to phantom accretion

$$t_{\text{ph}} = 2/3H_0 (16/27\rho_{\text{ph}}^0/\rho_{\text{rad}}^0)^{6-9/2(1+w)}$$



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- Phantom accretion to Hawking evaporation

$$m_t = \left[\frac{c^3}{G^2} \frac{A(m)}{|1+w|\rho_{\text{ph}}} \right]^{1/4}$$



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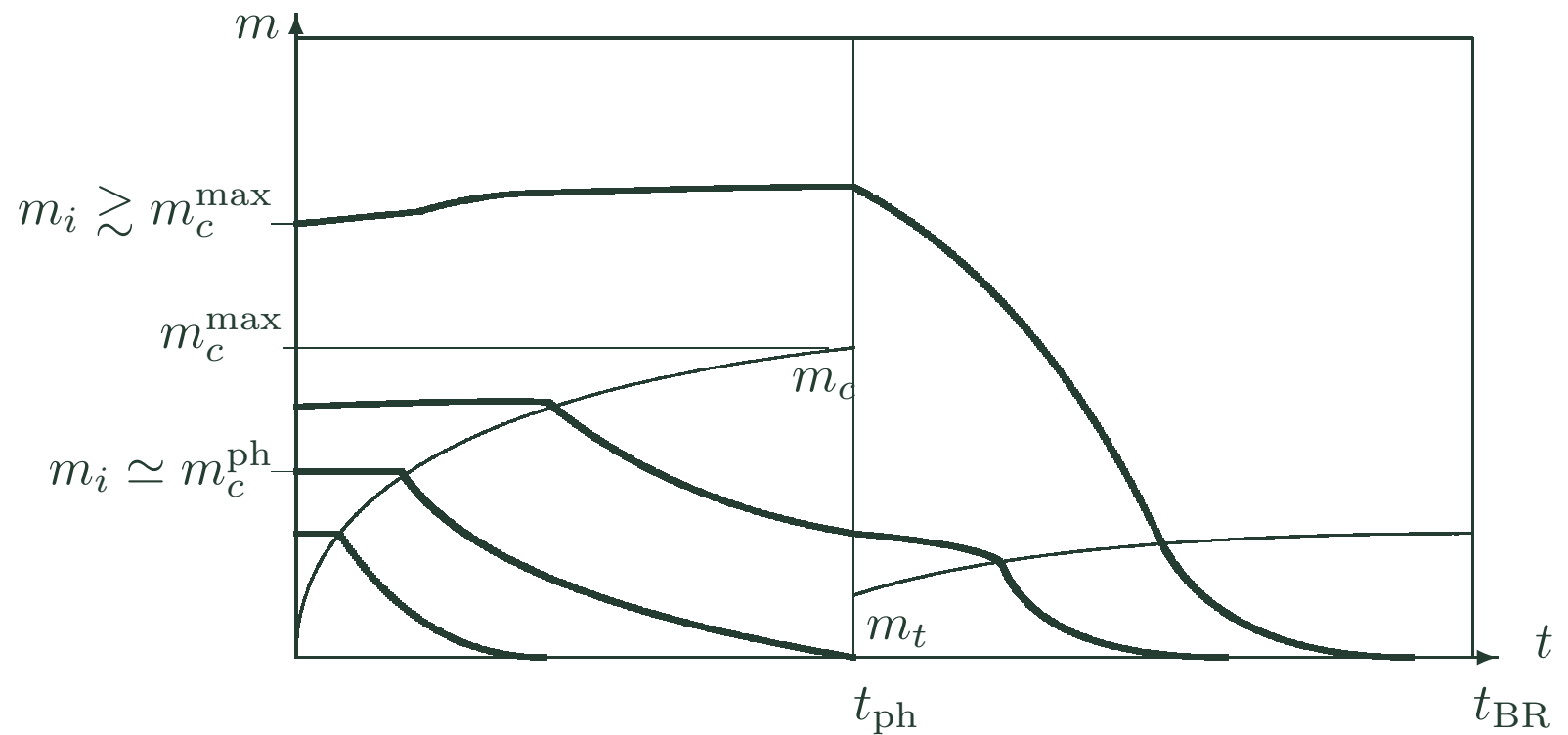
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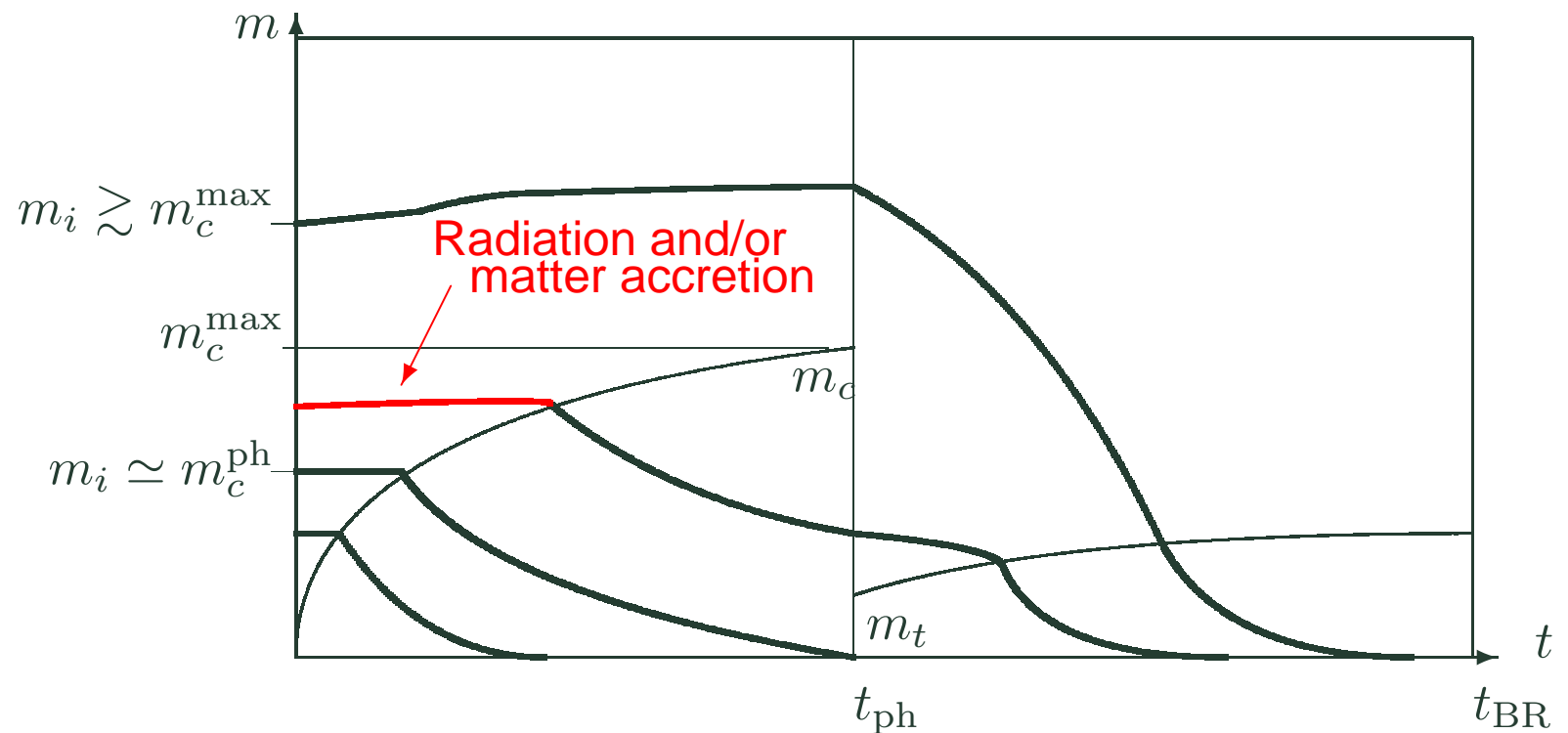
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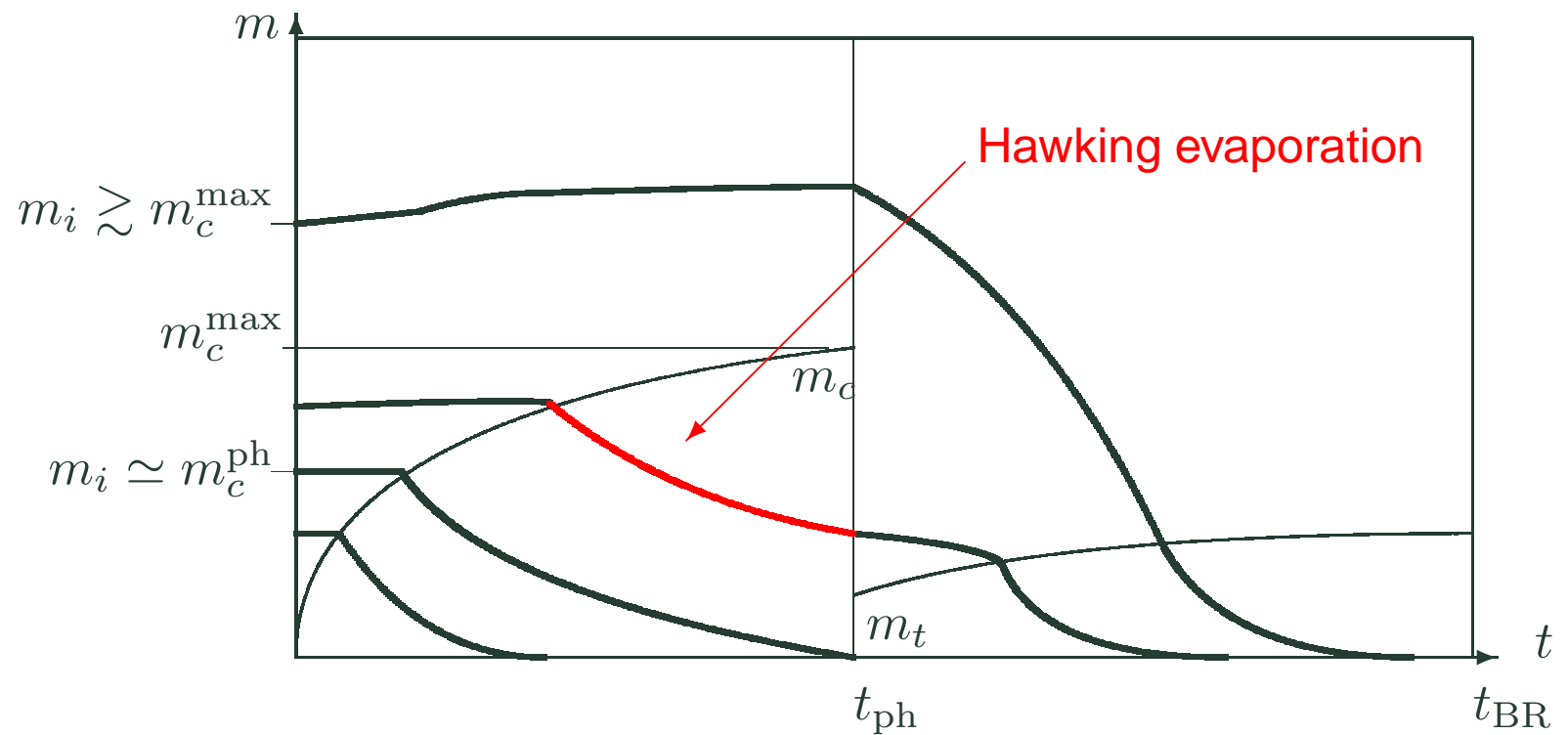
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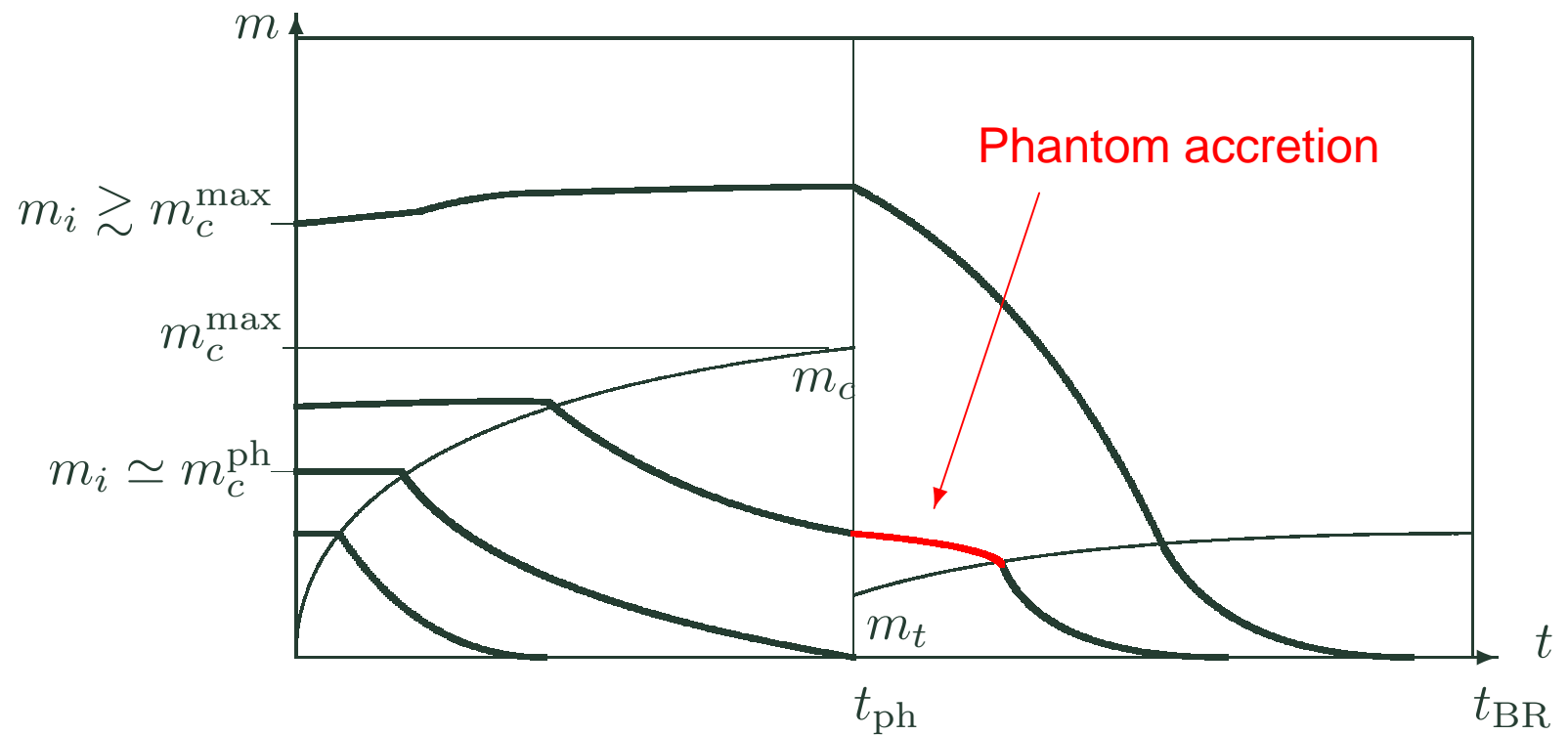
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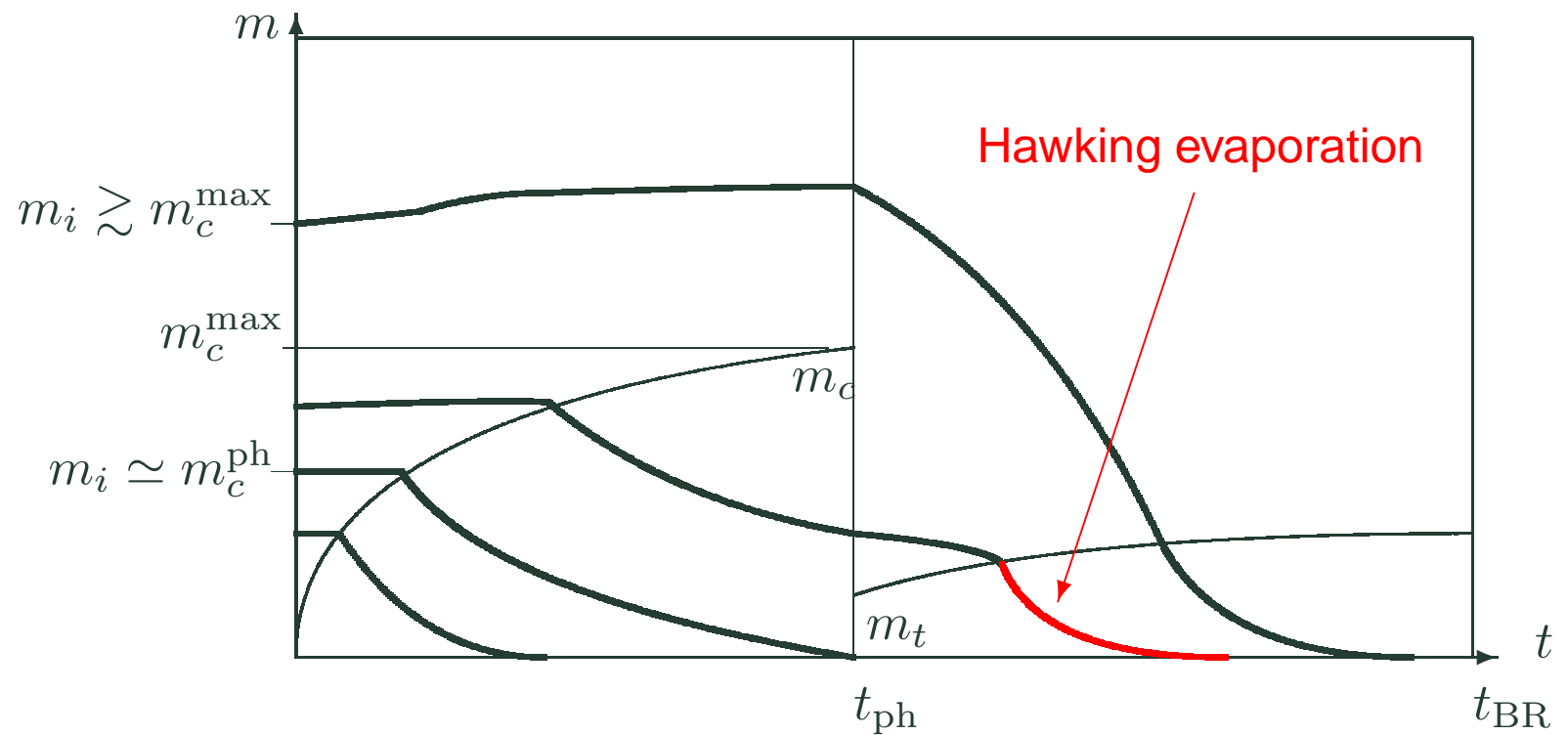
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- Simplest non-vacuum solution: Vaidya metric

$$ds^2 = (1 - 2m(\nu)/r) d\nu^2 - 2d\nu dr - r^2 d\Omega^2$$



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- Einstein equations only support null dust ($p = 0$)



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$$g_{00} = \left(1 - \frac{2m(\nu)}{r} + \lambda(\nu)r^2 + e(\nu)/r^2 \right)$$



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$$g_{00} = \left(1 - \frac{2m(\nu)}{r} + \lambda(\nu)r^2 + e(\nu)/r^2 \right)$$

Note: A red arrow points from the text "SdS" to the $\lambda(\nu)r^2$ term in the equation above.



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$$g_{00} = \left(1 - \frac{2m(\nu)}{r} + \lambda(\nu)r^2 + \cancel{e(\nu)/r^2} \right) \text{null dust}$$



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$$g_{00} = \left(1 - \frac{2m(\nu)}{r} + \lambda(\nu)r^2 + e(\nu)/r^2 \right)$$

- Cosmological black holes: generalized McVittie solution [Faraoni and Jacques, 2007]

$$ds^2 = \frac{\left(1 - \frac{m(t)}{2a(t)r}\right)^2}{\left(1 + \frac{m(t)}{2a(t)r}\right)^2} dt^2 - a^2(t) \left(1 + \frac{m(t)}{2a(t)r}\right)^4 (dr^2 + r^2 d\Omega^2)$$



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- Proposed fluids
 - Single perfect fluid

$$p = -\rho \quad \text{Schwarzschild-de Sitter}$$



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- Single perfect fluid

$$p = -\rho \quad \text{Schwarzschild-de Sitter}$$

- Imperfect fluid with heat transport
[Faraoni and Jacques, 2007, Gao et al., 2008]

$$\frac{dm}{dt} = \pm \frac{S(\rho + p)au^1 \left(1 - \frac{m}{2ar}\right)}{2} \sqrt{1 + a^2 \left(1 + \frac{m}{2ar}\right) (u^1)^2}$$



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- Two perfect dust-like fluids [Sultana and Dyer, 2005]

$$\frac{m(t)}{a(t)} = m_0$$



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$$T^{\mu\nu} = (\rho_1 + p_1) k^\mu k^\nu - p_1 g^{\mu\nu} + \rho_2 u^\mu u^\nu$$



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$$T^{\mu\nu} = \underbrace{(\rho_1 + p_1) k^\mu k^\nu - p_1 g^{\mu\nu}}_{\text{null dark energy}} + \rho_2 u^\mu u^\nu$$



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$$T^{\mu\nu} = (\rho_1 + p_1) k^\mu k^\nu - p_1 g^{\mu\nu} + \underbrace{\rho_2 u^\mu u^\nu}_{\text{dark matter}}$$



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- Einstein equations only allow *phantom* dark energy

$$(\rho_1 + p_1)(k^1)^2 + \rho_2(u^1)^2 = 0$$



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- Einstein equations only allow *phantom* dark energy

$$(\rho_1 + p_1)(k^1)^2 + \rho_2(u^1)^2 = 0$$

- Black hole mass evolution

$$\frac{dm}{dt} = -S \left(\frac{1 - \frac{m}{2ar}}{1 + \frac{m}{2ar}} \right)^4 \left[(\rho_1 + p_1) (k^0)^2 + \rho_2 \sqrt{ (k^0)^4 \frac{(\rho_1 + p_1)^2}{\rho_2^2} - \frac{\left(1 + \frac{m}{2ar}\right)^2}{\left(1 - \frac{m}{2ar}\right)^2} (k^0)^2 \frac{(\rho_1 + p_1)}{\rho_2} } \right] \quad (3)$$



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- Full solution to the Einstein equations must determine



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- Full solution to the Einstein equations must determine

- $\rho(r, t), p(\rho)$



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- $a(t)$



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- $u^\mu(r, t), k^\mu(r, t)$



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■ Full solution to the Einstein equations must determine

- $\rho(r, t), p(\rho)$
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 - $u^\mu(r, t), k^\mu(r, t)$
- } Already known for $r \rightarrow \infty$



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■ Approximate solutions might provide a simple back-reaction framework



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 - $m(t)$
- } Already known for $r \rightarrow \infty$

- Approximate solutions might provide a simple back-reaction framework
- Accretion of different types of fluids may be investigated
[Barrow, 1988]



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 - $a(t)$
 - $u^\mu(r, t), k^\mu(r, t)$
 - $m(t)$
- } Already known for $r \rightarrow \infty$

- Approximate solutions might provide a simple back-reaction framework
- Accretion of different types of fluids may be investigated
[Barrow, 1988]
- **Work in progress**



Introduction

Accretion models with no back reaction

Back Reaction analysis

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