# Constraining the Dark Energy Potential with BAO Surveys 

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## The Accelerating Universe: SNIa



## The Accelerating Universe



## The Friedmann Equations

Einstein's Equations + Homogeneity and Isotropy

$$
R_{\mu \nu}-\frac{1}{2} R g_{\mu \nu}=-\kappa T_{\mu \nu}
$$



FRW Metric with $\quad \kappa=8 \pi G \quad d s^{2}=d t^{2}-a^{2}(t)\left(\frac{d r^{2}}{1-k r^{2}}+r^{2} d \theta^{2}+r^{2} \sin ^{2} \theta d \phi^{2}\right)$
$H^{2}=\left(\frac{\dot{a}}{a}\right)^{2}=\frac{\kappa}{3} \rho-k \frac{c^{2}}{a^{2}} \quad \frac{\ddot{a}}{a}=-\frac{\kappa}{6}(\rho+3 p)$
For $\ddot{a}>0$ We need $3 p<-\rho$ but $\left\{\begin{array}{l}p=0 \text { for matter } \\ p=\frac{\rho}{3} \text { for radiation }\end{array}\right.$

## Cosmological Constant

Adding a Cosmological Constant $\Lambda$ to the Einstein Equation

$$
R_{\mu \nu}-\frac{1}{2} R g_{\mu \nu}-\Lambda g_{\mu \nu}=-\kappa T_{\mu \nu}
$$

The Friedmann Equations become

$$
\left\{\begin{array}{l}
H^{2}=\left(\frac{\dot{a}}{a}\right)^{2}=\frac{\kappa}{3} \rho+\frac{\Lambda}{3}-k \frac{c^{2}}{a^{2}} \quad \text { with } \quad \kappa=8 \pi G \\
\frac{\ddot{a}}{a}=-\frac{\kappa}{6}(\rho+3 p)+\frac{\Lambda}{3} \quad \text { If } \Lambda \text { dominates } \ddot{a}>0
\end{array}\right.
$$

## Modifying the Cosmological Constant

But what if $\Lambda$ is not a constant?
Alternative parameterizations of Dark Energy

- Constant equation of state

$$
p=-\rho=-\frac{\Lambda}{8 \pi G} \quad \Rightarrow \quad p=w \rho
$$

- More general equations of state

$$
\begin{gathered}
w(z)=w_{0}+w^{\prime} z \quad w(a)=w_{0}+w_{a}(1-a) \\
\text { DE task force }
\end{gathered}
$$

## A More general parameterization

But nature could be much more general than that...
If Dark Energy came from the potential of a scalar field $V(\phi)$
$\rho_{\phi}=K+V(\phi) \quad p_{\phi}=K-V(\phi) \quad$ with $\quad K=\frac{1}{2} \dot{\phi}^{2}$

The Friedmann Equations become

$$
\left\{\begin{array}{l}
H^{2}=\frac{\kappa}{3}(\rho+V+K)-k \frac{c^{2}}{a^{2}} \\
\frac{\ddot{a}}{a}=-\frac{\kappa}{6}(\rho+3 p+4 K-2 V(\phi))
\end{array}\right.
$$

If $V(\phi)$ dominates we can also have $\ddot{a}>0$
but $V(z)$ could be a very general function of $z$

## A More general parameterization

Even if $V(z)$ is a general function we need to parameterize it and try to fit it to data

We expand the potential in Chebyshev polynomials

$$
V(z) \approx \sum_{n=0}^{N} \lambda_{n} T_{n}\left(2 \frac{z}{z_{\max }}-1\right)
$$

with $z_{\text {max }}$ the maximum
redshift of the survey
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We will study the constraints that BAO surveys can place on the first three coefficients $\lambda_{i}$

## BAO



## Dark Matter Baryons Photons Neutrinos

Animation by D. Eisenstein

## BAO



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## BAO



Dark Matter Baryons Photons Neutrinos

Animation by D. Eisenstein

## BAO

The distance traveled by sound until recombination

$$
\begin{aligned}
\text { Rs }= & 153.3 \pm 2.0 \text { Mpc WMAP 5th year } \\
& \text { provides a "standard ruler" }
\end{aligned}
$$

Measuring its transverse size $\Delta \theta$ gives:

$$
d_{A}^{c o}=\int_{0}^{z} \frac{c}{H\left(z^{\prime}\right)} d z^{\prime}=\frac{R_{s}}{\Delta \theta}
$$

Measuring its radial size $\Delta z$ gives:

$$
H(z)=\frac{c \Delta z}{R_{s}}
$$

## PAU (Physics of the Accelating Universe)

PAU is a photometric survey but with $\sim 40$ filters of $\sim 100 \AA$
Delta $\lambda=100 \AA$



Figure by T. Benítez
This allows to measure the redshift with $\sigma_{z} \sim 0.0035(1+z)$
Will measure between $\mathrm{z}=0.1-1$

## Adept (Advanced DE Physics Telescope)

Space-based telescope to probe DE through observation of SN Ia and BAO

Spectroscopic redshifts


Can go deeper in redshift, will measure between $z=1-2$

## BAO Errors

In C. Blake et al. 2005 several BAO surveys were simulated and a formula to estimate the precision of the survey was fitted

$$
\begin{array}{rlr}
\sigma_{H}\left(z_{i}\right) & =x_{0}^{H} \frac{4}{3} \sqrt{\frac{V_{0}}{V_{i}}} f_{n l}\left(z_{i}\right) & \sigma_{d}\left(z_{i}\right)=x_{0}^{d} \frac{4}{3} \sqrt{\frac{V_{0}}{V_{i}}} f_{n l}\left(z_{i}\right) \\
f_{n l}\left(z_{i}\right) & =\left\{\begin{array}{ccc}
1 & z>1.4 & x_{0}^{d}=0.0085 \\
\left(\frac{1.4}{z}\right)^{1 / 2} & z<1.4 & x_{0}^{H}=0.0148
\end{array}\right. \\
\hline
\end{array}
$$

We have studied two future BAO surveys: PAU and Adept

| PAU | Adept |
| :---: | :---: |
| $A=10000$ square degrees | $\mathrm{A}=30000$ square degrees |
| 9 redshift bins | 10 redshift bins |
| between 0.1 and 1 | between 1 and 2 |

## The importance of measuring along the line of sight

Using BAO as Standard Ruler
Information on $d_{A}(z)$ only
Information on $\mathrm{H}(\mathrm{z})$ only



Adept 1, 2 and $3 \sigma$ constraints on the first 3 coefficients Assuming a $\Lambda$ CDM model $\Omega_{\text {m0 }}=0.24$ and $1 \sigma$ priors on $\Omega_{\mathrm{m} 0} \mathrm{~h}^{2}(0.01), \Omega_{\mathrm{k0}}(0.03)$ and $\mathrm{H}_{0}(8 \mathrm{Km} / \mathrm{s} / \mathrm{Mpc})$

## PAU

Using BAO as Standard Ruler


PAU 1, 2 and $3 \sigma$ constraints on the first 3 coefficients
Assuming a $\Lambda$ CDM model $\Omega_{\mathrm{mo}}=0.24$
and $1 \sigma$ priors on $\Omega_{m 0} h^{2}(0.01), \Omega_{k 0}(0.03)$ and $\mathrm{H}_{0}(8 \mathrm{Km} / \mathrm{s} / \mathrm{Mpc})$

## Adept

Using BAO as Standard Ruler



Adept 1, 2 and $3 \sigma$ constraints on the first 3 coefficients Assuming a $\Lambda$ CDM model $\Omega_{\mathrm{m} 0}=0.24$ and $1 \sigma$ priors on $\Omega_{m 0} h^{2}(0.01), \Omega_{\mathrm{k} 0}(0.03)$ and $\mathrm{H}_{0}(8 \mathrm{Km} / \mathrm{s} / \mathrm{Mpc})$

## Constraints on the potential



1 and $2 \sigma$ constraints on the potential

## Constraints on the potential


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1 and $2 \sigma$ constraints on the potential

## DE Equation of State

What about parameterizing Dark Energy through its equation of state $w(z)$ ?
$w(z)$ can also be expanded in Chebyshev Polynomials:

$$
w(z) \approx \sum_{n=0}^{N} w_{n} T_{n}\left(2 \frac{z}{z_{\max }}-1\right)
$$

If $w>-1$ the two descriptions are equivalent

## Constraints on the equation of state



1 and $2 \sigma$ constraints on the equation of state

## Conclusions

- A non parametric reconstruction of dynamical DE requires a precise measurement of $H(z)$ and $H^{\prime}(z)$
- BAO can directly probe $\mathrm{H}(\mathrm{z})$
- $\mathrm{H}^{\prime}(\mathrm{z})$ much more challenging observationally
- Lacking a measurement of $\mathrm{H}^{\prime}(z)$ we need a parameterization of DE to fit data
- Future BAO experiments can constrain $\mathrm{V}(\mathrm{z})$ and $\mathrm{w}(\mathrm{z})$
- Excellent probes of dynamical DE


## Constraints on the potential



1 and $2 \sigma$ constraints on the potential

## Reconstruction of the potential

From the Friedmann Equations

$$
\begin{aligned}
& \qquad V(z)=\left(3-\varepsilon_{1}\right) \frac{H^{2}}{\kappa}+\frac{1}{2}\left(p_{T}-\rho_{T}\right) \\
& \text { with } \quad \varepsilon_{1}=-\frac{\dot{H}}{H^{2}}=1-\frac{\ddot{a}}{a} H^{-2}=\frac{d H}{d z} \frac{(1+z)}{H}
\end{aligned}
$$

$H(z)$ and $H^{\prime}(z)$ needed to reconstruct $V(z)$

## Galaxy Ages

## Using Relative Galaxy Ages as Standard Chronometers



1 and $2 \sigma$ contours
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## Supernovae

## Using SN as Standard Candles




1 and $2 \sigma$ contours
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From the supernovae data of A. G. Riess et al astro-ph/0402512

