

Cosmic Superstrings & their Astrophysical Consequences

Anastasios Avgoustidis
(ECM & ICC Barcelona)

work with E.P.S. Shellard
(DAMTP, Cambridge)

[[hep-ph/0410349](#), [hep-ph/0504049](#),
[astro-ph/0512582](#), [Arxiv:0705.3395](#)]

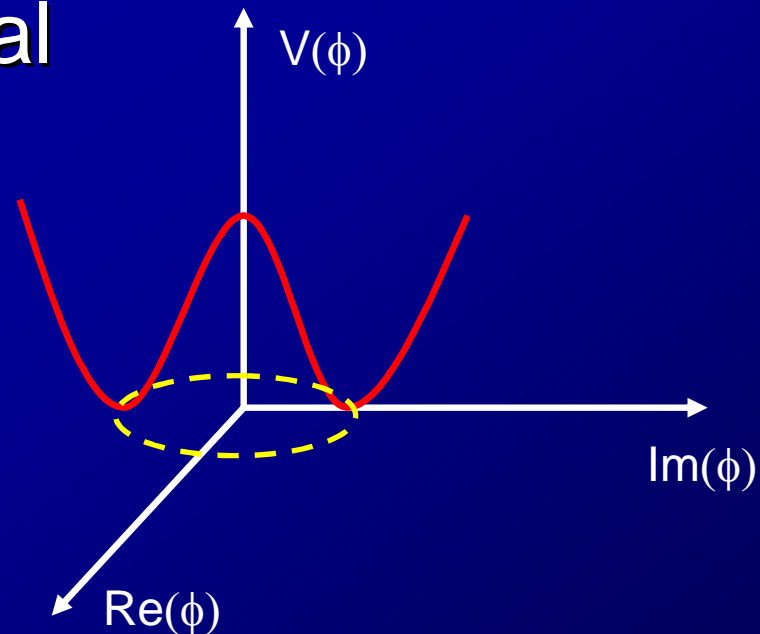


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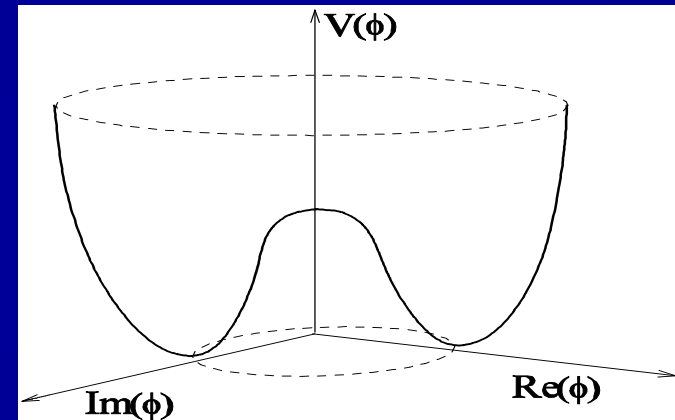
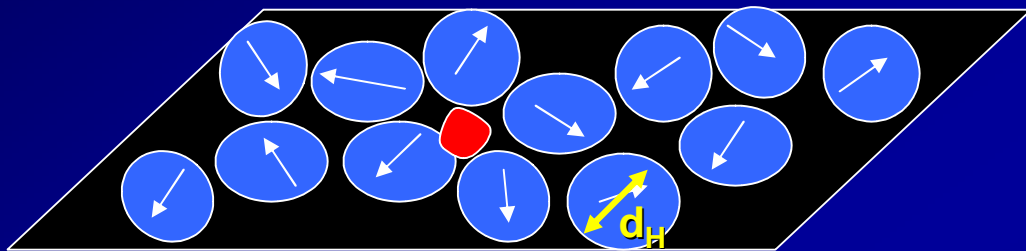
Cosmic Strings

- Cosmic strings: Line-like concentrations of energy arising as topological defects in cosmological phase transitions.
- *Example*: Complex scalar field with mexican hat potential



Kibble Mechanism: 2D Example

Vortex Formation

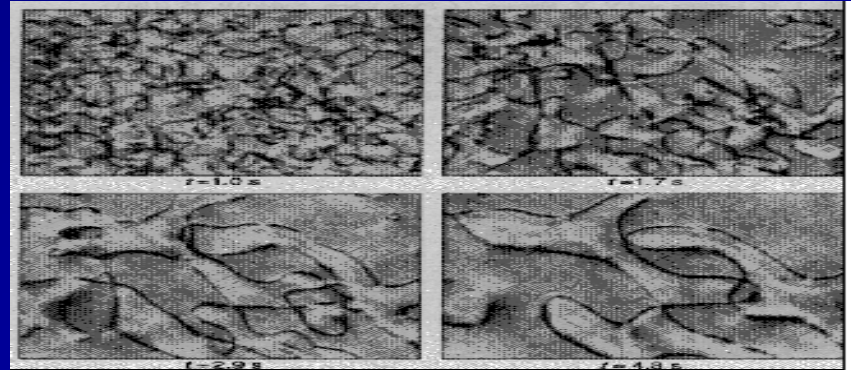


Kibble Mechanism: Any topologically allowed defect will be cosmologically formed

Why Strings?

■ Produced in Phase Transitions

■ Until 1997:
Structure Formation



$$G\mu \approx 10^{-6} \approx \left(\frac{\Delta T}{T} \right)_{CMB}$$

Ruled out (Battye et al)

■ Now:

String Theory, Brane Inflation

(Quevedo et al 2001,
Sarangi & Tye 2002)

Generically produced in SUSY GUTS

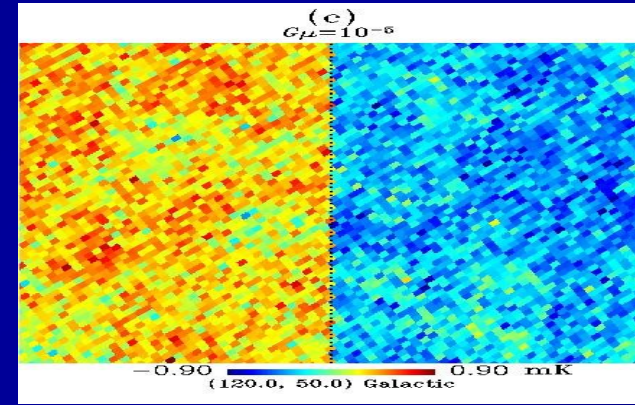
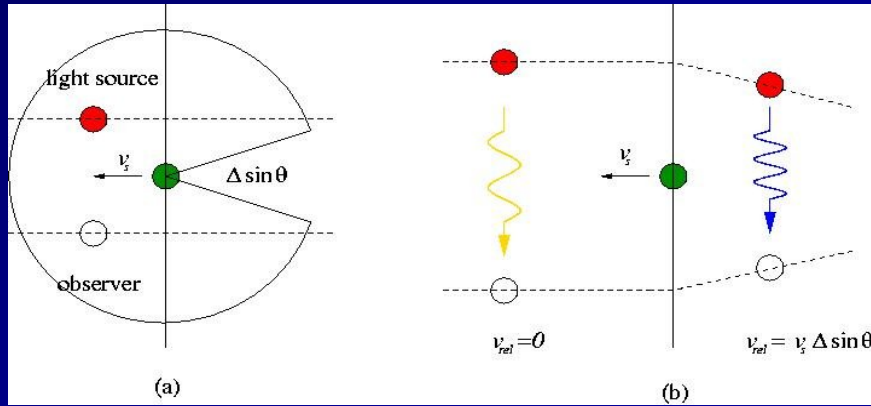
(Jeanerot et al 2003)

Observational: ~~CBL1~~, Oscillating Loop
(Arguably) favoured by CMB

(Bevis, Hindmarsh, Kunz & Urrestilla, astro-ph/0702223)

String Gravity & Cusps

■ String Gravity: (Vilenkin 1981, Gott 1985)



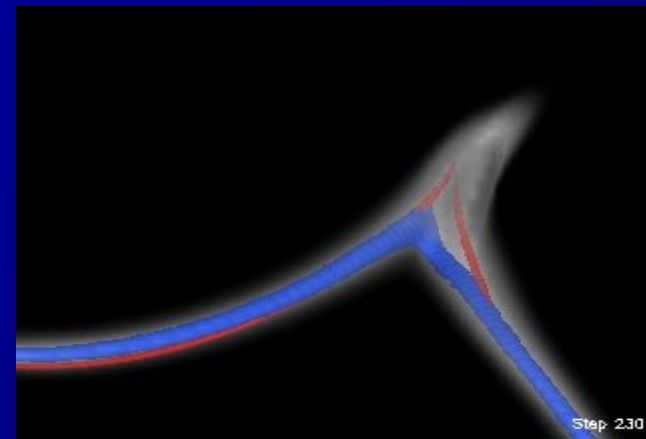
(Kaiser & Stebbins 1984)

■ Radiation from Cusps:

(Brandenberger 1987
Bhattacharjee 1989
Damour & Vilenkin 2001)

Astrophysical Consequences

(Berezinsky et al 1998, Vachaspati 2008)



(Potential) Observational Effects

- CMB:
 - a) Discontinuous Doppler Shift by strings at “present” time (Kaiser-Stebbins effect)
 - b) Fluctuations at surface of last scattering
 - c) Sachs-Wolf effect, gravitational waves,...

B-mode polarization (Pogosian et al 2006, 2007)
No degeneracy with primordial tensors (Urethilla et al 2008)
- Gravitational Lensing: Double Images (no distortion)
Microlensing (Kuijken et al 2007)
- Gravitational Radiation:
 - a) Stochastic Background from loops (pulsar timing)
 - b) Gravitational Waves from Cusps (LIGO, LISA)

(Damour & Vilenkin 2000, 2001, 2005)

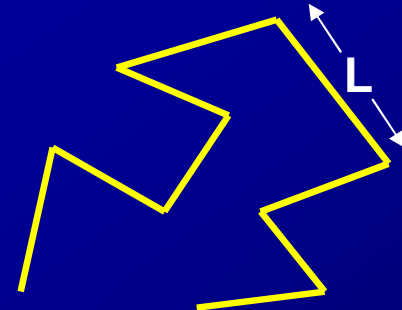
String Evolution

■ Field Theory Simulations

■ Nambu Strings $\begin{cases} \rightarrow \text{Simulations} \\ \rightarrow \text{Macroscopic approach} \end{cases}$

String network is Brownian

Characterised by correlation length L



Define energy density

$$\rho = \frac{\mu L}{L^3} = \frac{\mu}{L^2}$$

$$\dot{\rho} \approx -2H\rho - \rho/L$$

loop production

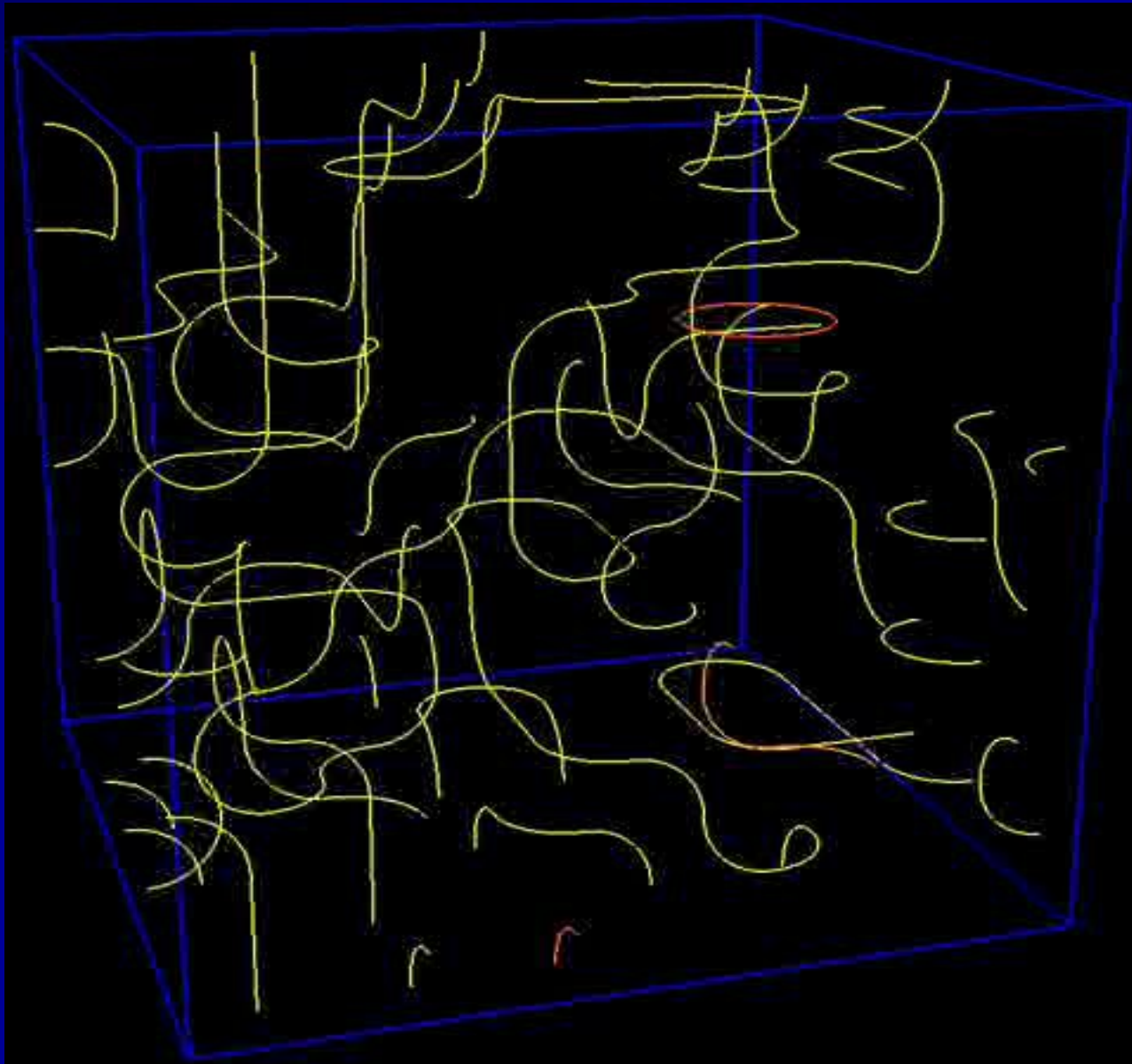
Find scaling solution

$$L \sim t$$

(Kibble 1985)

Nambu-Goto Simulation

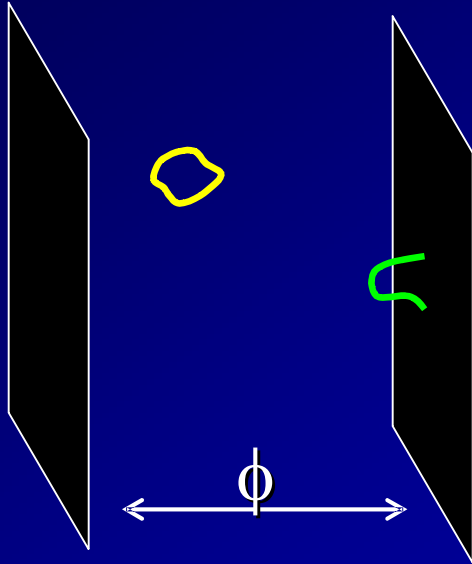
(Martins & Shellard)



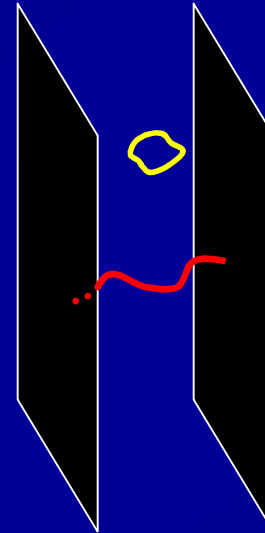
Analytic Models

- Improved One-Scale Model (Bennett 1986)
- Kink-Counting Model (Allen & Caldwell 1990; Austin 1993)
- Functional Approach (Embacher 1992)
- Three-Scale Model (Austin, Copeland & Kibble 1993)
- Wiggly Model (Martins 1997)
- Velocity Dependent One-Scale (VOS) Model
(Martins & Shellard 1996/2000)

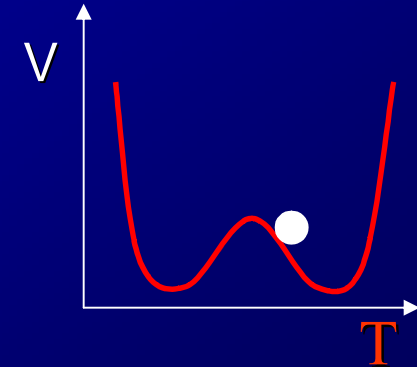
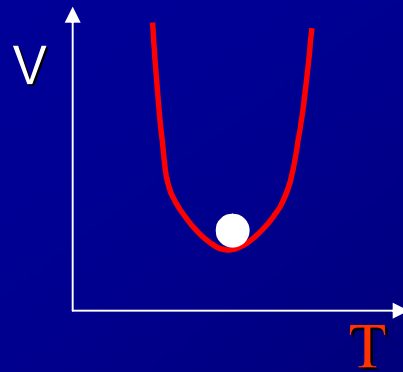
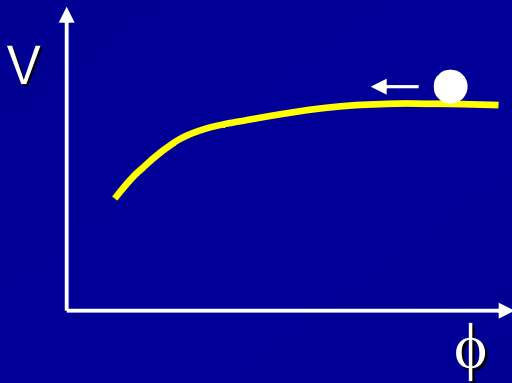
Brane Inflation



$$V_{\text{int}}(\phi) \approx -\frac{c}{\phi^{d-2}}$$



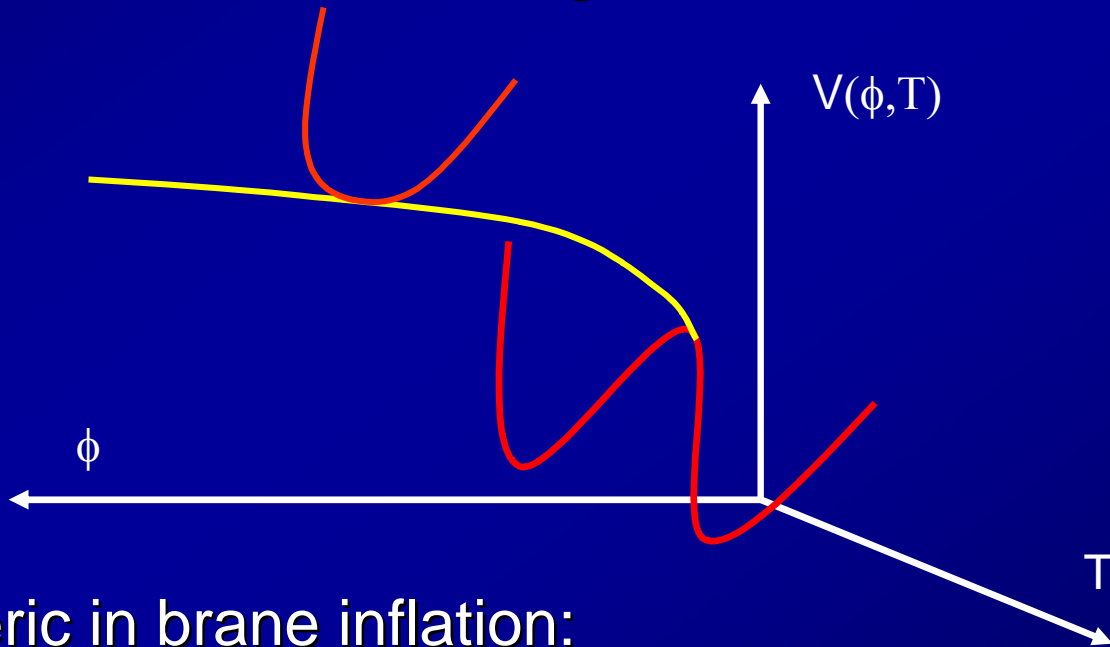
$$M^2(\phi) \approx M_s^2 \left(\frac{\phi}{\phi_c} - 1 \right)$$



Relic String Defects

Brane inflation ends with tachyon condensation, producing a network of F & D strings.

(BMNQRZ 2001,
Sarangi & Tye 2002,
Dvali & Vilenkin 2004)



This seems to be generic in brane inflation:

- Brane-Antibrane (BMNQRZ 2001, KKLMNT, 2004)
- Branes at angles (Garcia-Bellido, Rabadan & Zamora 2001, Gomez-Reino & Zavala 2004)
- D3/D7 (Dasgupta, Herdeiro, Hirano, Kallosh 2002)
- Wilson Line Inflation (AA, Cremades & Quevedo 2004)

Strings in Extra Dimensions

Strings can miss in $D > 3$ spatial dimensions

Introduces a probability for intercommuting $P < 1$
(Jones, Stoica & Tye 2003)

$$\dot{\rho} \approx -2H\rho - \rho P/L$$

Scaling Solution:

$$\rho \propto P^{-2}$$

- Would lead to a much denser string network
- $P < 1$ also enhances gravitational radiation
(Damour & Vilenkin 2005)

The EDVOS Model 1: Microscopic Equations

■ Metric $ds^2 = N(t)^2 dt^2 - a(t)^2 d\mathbf{x}^2 - b(t)^2 d\mathbf{l}^2$

■ Nambu-Goto Action $S = -\mu \int \sqrt{-\gamma} d^2\sigma$

■ Equations of motion $\dot{\varepsilon} = -N^{-2} \varepsilon \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\varepsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\varepsilon} \right)^2 \right] \right\}$

$$\ddot{\mathbf{x}} + \left\{ 2 \frac{\dot{a}}{a} - N^{-2} \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\varepsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\varepsilon} \right)^2 \right] \right\} \right\} \dot{\mathbf{x}} = \left(\frac{\mathbf{x}'}{\varepsilon} \right) \varepsilon^{-1}$$

$$\ddot{\mathbf{l}} + \left\{ 2 \frac{\dot{b}}{b} - N^{-2} \left\{ N\dot{N} + a\dot{a} \left[\dot{\mathbf{x}}^2 - \left(\frac{\mathbf{x}'}{\varepsilon} \right)^2 \right] + b\dot{b} \left[\dot{\mathbf{l}}^2 - \left(\frac{\mathbf{l}'}{\varepsilon} \right)^2 \right] \right\} \right\} \dot{\mathbf{l}} = \left(\frac{\mathbf{l}'}{\varepsilon} \right) \varepsilon^{-1}$$

■ Energy-Momentum Tensor & Energy

$$T^{\mu\nu} = \frac{1}{Na^3b^{D-3}} \mu \int d\sigma \left(\varepsilon \dot{x}^\mu \dot{x}^\nu - \varepsilon^{-1} x'^\mu x'^\nu \right) \delta^{(D)}(\mathbf{x} - \mathbf{x}(\sigma, t), \mathbf{l} - \mathbf{l}(\sigma, t))$$

$$E(t) = \int_{t=const} \sqrt{h} n_\mu n_\nu T^{\mu\nu} d^3\mathbf{x} d^{D-3}\mathbf{l} = N(t) \mu \int \varepsilon d\sigma$$

The EDVOS Model 2: Macroscopic Equations

- Energy Density (equiv. correlation length L)

$$2 \frac{dL}{dt} = \left[(2 + w_l^2) + (2 - w_l^2) v_x^2 + (1 - w_l^2) v_l^2 \right] HL + cP_{eff} v_x$$

Modify 3D terms

New term due to
extra dimensional
velocities

- String Velocities

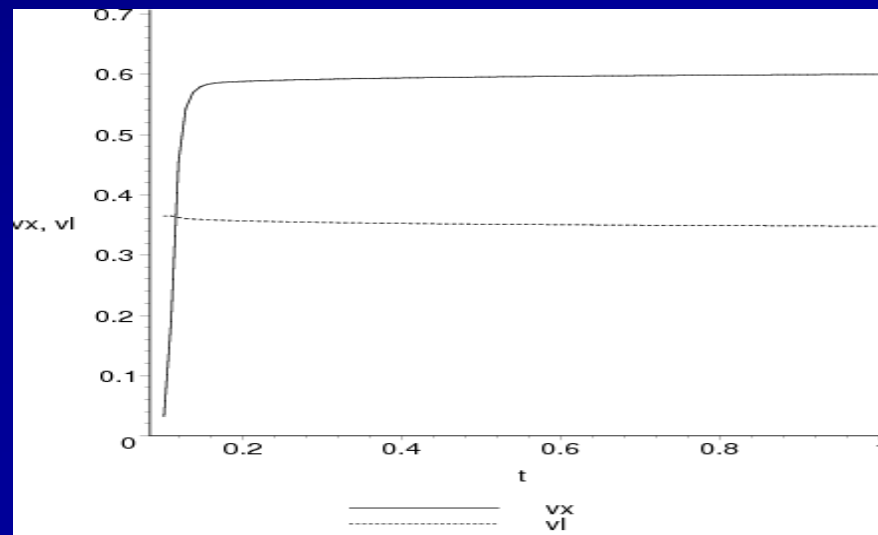
$$v_x \frac{dv_x}{dt} = \frac{k_x v_x}{L} (1 - v^2) - (2 - w_l^2) H v_x^2 (1 - v^2) - H v_x^2 v_l^2$$
$$v_l \frac{dv_l}{dt} = \frac{k_l v_l}{L} (1 - v^2) - (1 - w_l^2) H v_l^2 (1 - v^2) + H v_l^2 v_x^2$$

Effective 3D string motion slows down due to extra dimensional velocities

(AA & Shellard 2004)

The EDVOS Model 2: Macroscopic Equations

$$2 \frac{dL}{dt} = \left[(2 + w_l^2) + (2 - w_l^2) v_x^2 + (1 - w_l^2) v_l^2 \right] HL + cP_{eff} v_x$$



$$v_x^2 + v_l^2 \leq 1/2$$

$$v_x \frac{dv_x}{dt} = \frac{k_x v_x}{L} (1 - v^2) - (2 - w_l^2) H v_x^2 (1 - v^2) - H v_x^2 v_l^2$$

$$v_l \frac{dv_l}{dt} = \frac{k_l v_l}{L} (1 - v^2) - (1 - w_l^2) H v_l^2 (1 - v^2) + H v_l^2 v_x^2$$

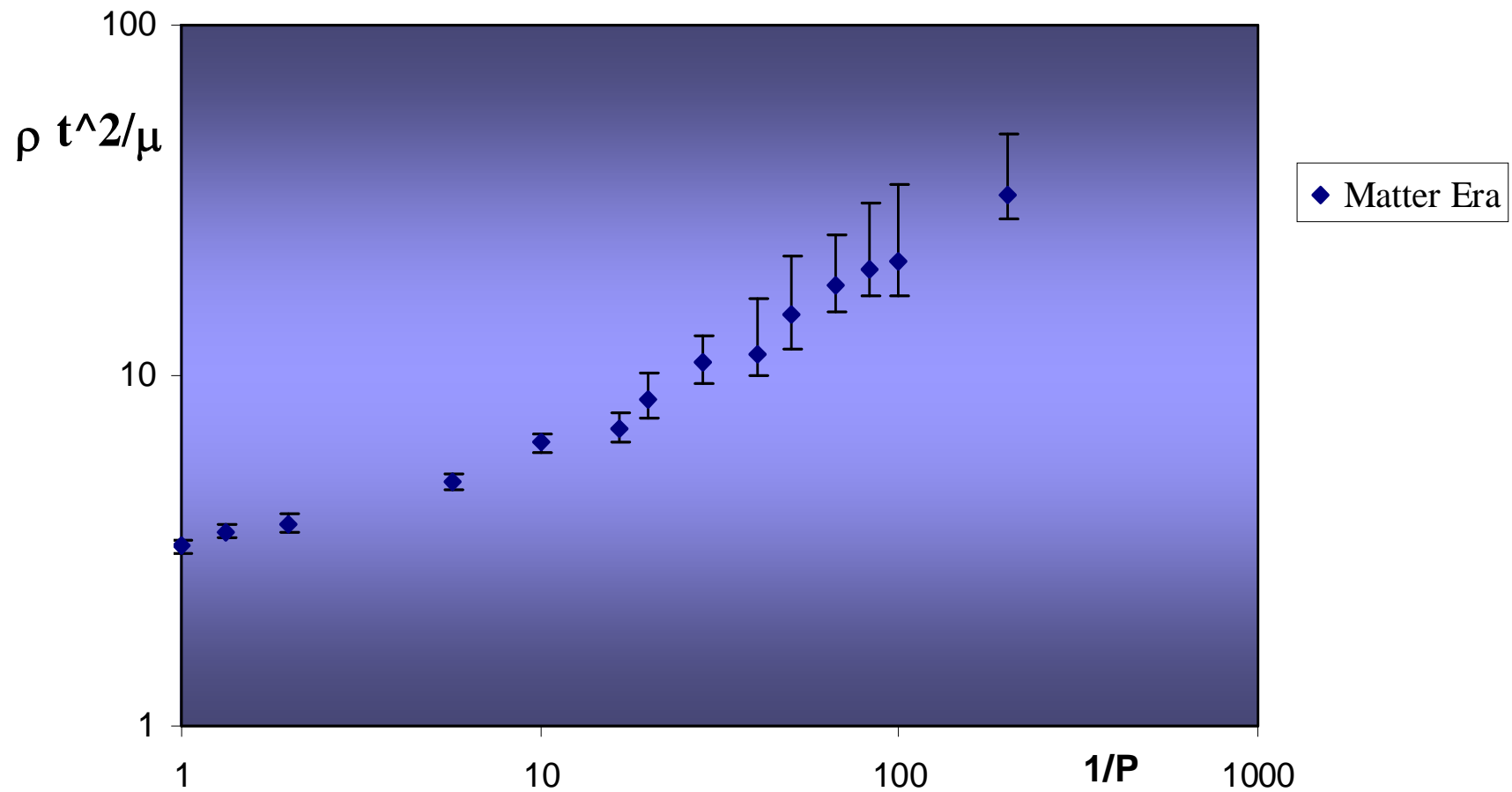
Intercommuting Probability

- Jackson, Jones & Polchinski 2004: $0.001 < P < 1$
+Kinematic Constraints (Copeland, Kibble & Steer 2006,
Salmi et al 2007)
- Effect of Small-Scale Structure. Need simulations.
- Sakellariadou 2004: Flat space simulations with
 $P < 1$ suggest $\rho \propto P^{-1}$
- Simulations in expanding space: $\rho \propto P^{-0.6 \pm 0.1}$
(AA & Shellard 2005)
Cf Vanchurin 2007 $\rho \propto P^{-2/3}$

Numerical Results

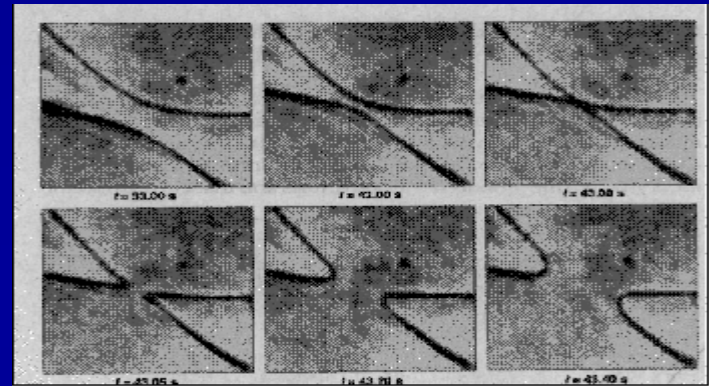
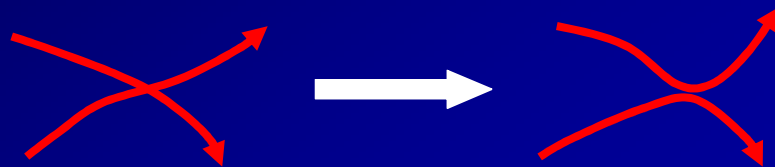
(AA & Shellard 2005)

String Density vs Inverse Intercommuting Probability

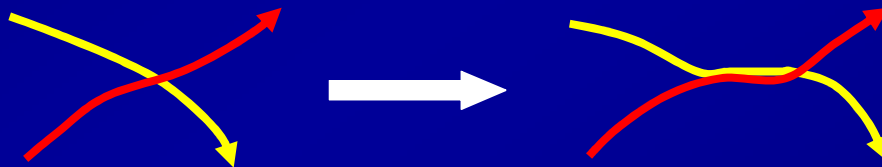


String Junctions

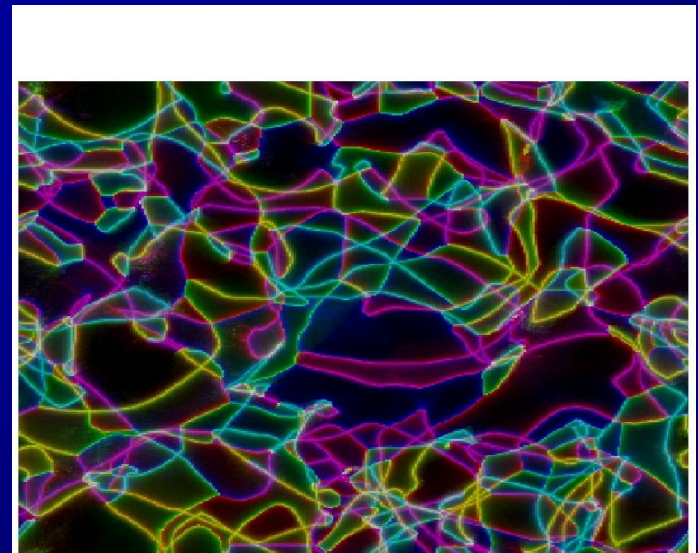
- Usual field theory strings interact by exchange of partners ($P \sim 1$)



- F and D strings can bind together to form F-D composites



Cf. Non-Abelian Strings



Scaling?

- **Copeland & Saffin, 2005**: Numerical evidence for scaling in field theory model.
Also **Hindmarsh & Saffin, 2006**
- **Tye, Wasserman & Wyman, 2005**: Multi-tension VOS model for entangled string network.
Evidence for scaling.
- **AA & Shellard, 2006**: Non-Abelian Velocity-Dependent One-Scale Model (NAVOS) for networks with junctions.

NAVOS: Macroscopic Equations

■ String Densities:

$$\dot{\rho}_i = -2 \frac{\dot{a}}{a} (1 - 2v_i^2) \rho_i - \frac{c_i v_i \rho_i}{L_i} - \sum_{a,k} \frac{d_{ia}^k \tilde{v}_{ia} l_{ia}^k(t)}{L_a^2 L_i^2} + \sum_{b,a \leq b} \frac{d_{ab}^i \tilde{v}_{ab} l_{ab}^i(t)}{L_a^2 L_b^2}$$

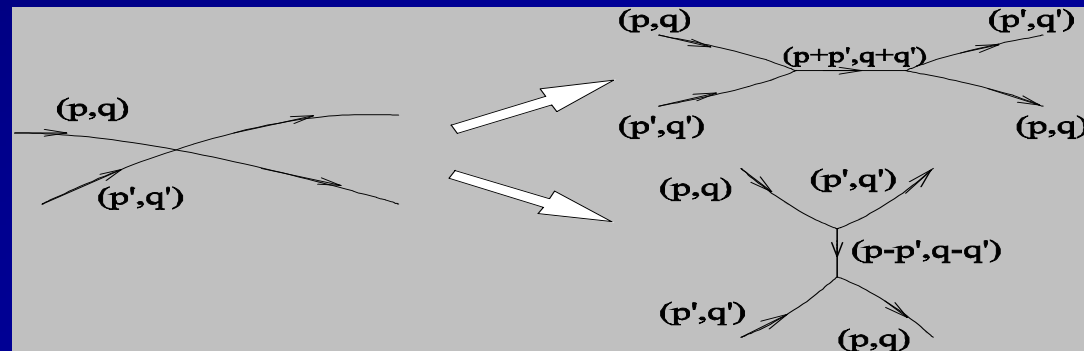
■ String Velocities:

$$\dot{v}_i = (1 - v_i^2) \left[\frac{k_i}{L_i} - 2 \frac{\dot{a}}{a} v_i + \sum_{b,a \leq b} d_{ab}^i \frac{\tilde{v}_{ab}}{v_i} \frac{(\mu_a + \mu_b - \mu_i)}{\mu_i} \frac{l_{ab}^i(t) L_i^2}{L_a^2 L_b^2} \right]$$

where $l_{ij}^k(t)$ are model dependent $f(L_i, v_i, \mu_i)$
and c_i, d_{ij}^k are free parameters.

Cosmic Superstrings

- For (p,q) -strings zipping interactions must be included:

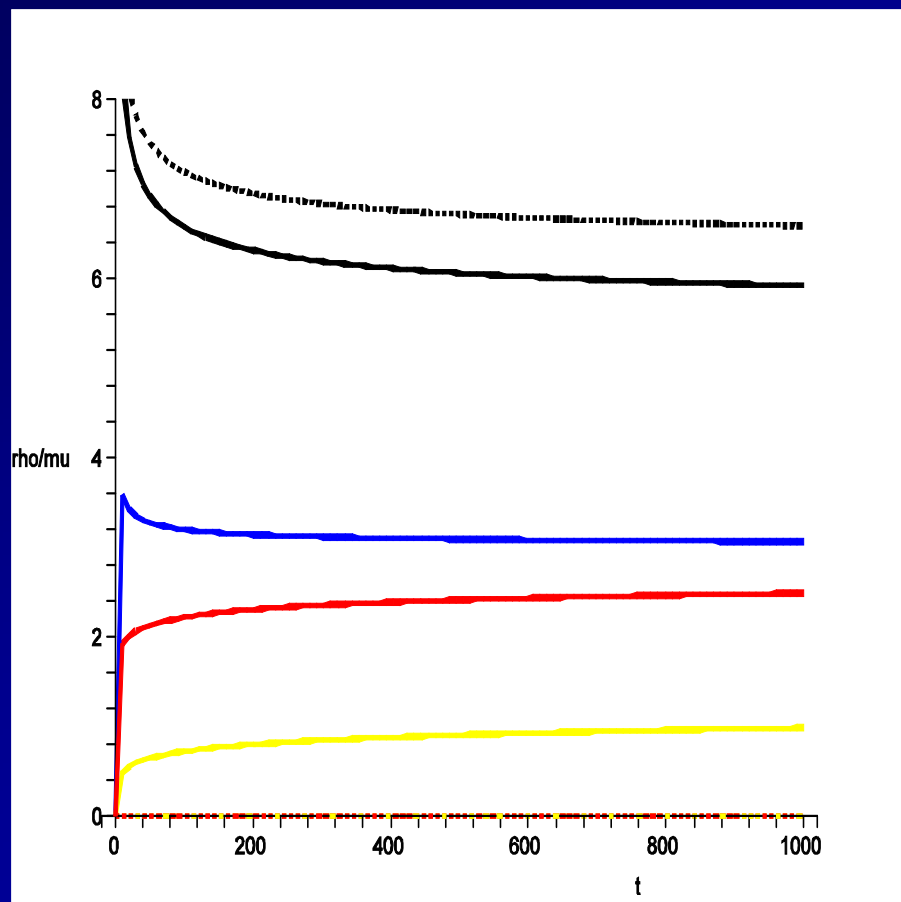


- Probability of additive/subtractive process is:

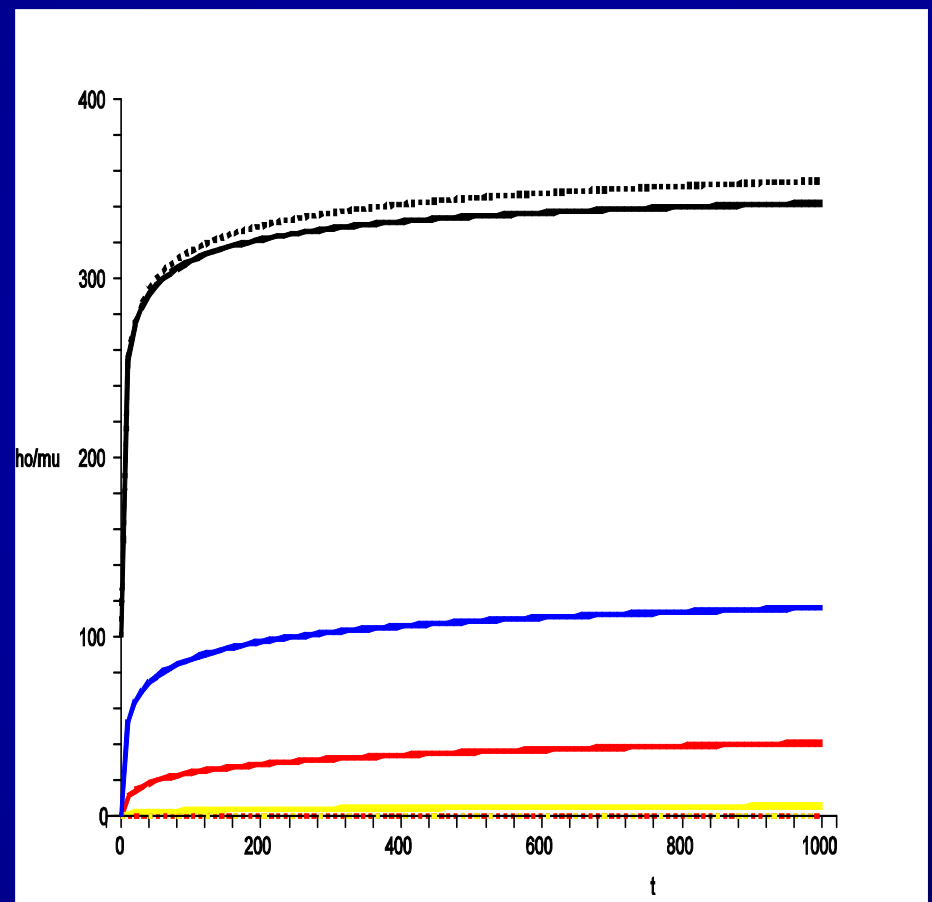
$$P_{(p,q)(p',q')}^{\pm} = \frac{1}{2} \left[1 \mp \left(\frac{pp'g_s^2 + qq'}{(p^2g_s^2 + q^2)^{1/2}(p'^2g_s^2 + q'^2)^{1/2}} \right) \right]$$

(Tye et al 2005)

Scaling Results



$P=1$



$P \ll 1$

Cosmological Constraints

Observational Bounds on Cosmic Strings:

$$G\mu < 6 \times 10^{-7} \quad \text{CMB}$$

(Wyman et al 2005)

$$G\mu < 5 \times 10^{-8} \quad \text{Pulsar Timing}$$

(Kaspi et al 1994, Mc Hugh et al 1996)

Impose constraints on stringy parameters:

String Scale, Compactification Radii, Warping, String coupling,...

DEGENERACIES

Degeneracies could be broken using other constraints:

CMB, Non-Gaussianity, Gravitational Waves

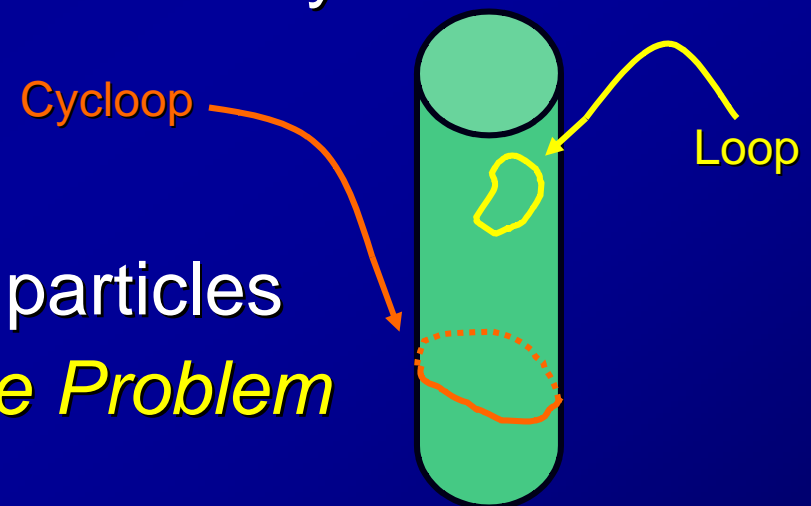
(Lorenz, Martin & Ringeval, 2007)

More Stringy Constraints

Example: Cycloops (AA & Shellard 2005)

If internal manifold admits non-trivial 1-cycles then loops can be trapped

Cycloops behave like massive particles from 3D point of view *Monopole Problem*



Imposes stringent constraints on energy scale of inflation

Other examples: Dilaton emission (Babichev et al 2005)

Vortons (Brax et al 2006)

Outlook & Open Questions

- Cosmic Strings provide a potential observational window into HEP
- Already constraining stringy parameters but need better quantitative understanding
- Some Open Problems:
 - Small-Scale Structure (Martins & Shellard 2005, Rocha & Polchinski 2006, Sakelariadou et al, Vanchurin et al)
 - Number of kinks in loops
 - Networks with Junctions: Scaling?
 - Non-Gaussianity