# starox: stellar evolution code

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Domain of applicability: Pre Main Sequence to Post Main Sequence

## 1. Basic structure of code

- 1. Mesh in mass M(i), i = 0, N; structure variables  $V_i(i)$ , Chemical Species  $X_k(i)$
- 2. At time t, predict values  $V_j, X_j$  at t + dt, calculate new  $X_k$  at t + dt
- 3. Solve structure for  $V_i$  at t + dt with new  $X_i$
- 4. Iterate steps 2 and 3 to convergence.  $V_j = r, L, \rho, T, P, U, \ldots; :X_k(k = 1, 9) = {}^{1}H, {}^{3}He, {}^{4}He, {}^{12}C, {}^{13}C, {}^{14}N, {}^{15}N, {}^{16}O, {}^{17}O$

## 2. Structure equations in form solved in code

$$\frac{\partial M_r}{\partial r^3} = \frac{4}{3}\pi G\rho, \qquad \frac{\partial L_r}{\partial M_r} = \epsilon - \left[\frac{\partial U}{\partial t} - \frac{P}{\rho^2}\frac{\partial \rho}{\partial t}\right], \quad \frac{\partial T}{\partial r^2} = -\nabla \frac{T}{P}\frac{GM_r\rho}{2r^3}$$
$$\nabla \equiv \frac{\partial \log T}{\partial \log P} = \nabla_{rad} \text{ (if } \leq \nabla_{ad} \text{) else} = \nabla_{con}. \quad \nabla_{rad} = \frac{3L_rP\kappa}{64\pi\sigma GM_rT^4}, \ \nabla_{con} \text{ from } MLT$$

MLT Convective Model as implemented in this code,  $\ell = \alpha H_p$ 

$$\nabla_{con} = \nabla_{ad} + \Delta \nabla, \quad \Delta \nabla = \left(\frac{2\,\rho B^2}{\lambda P}\right) \left(x^2 + x\right), \quad B = \frac{48\sigma T^3}{c_p \kappa \alpha^2 H \rho^2}, \quad \lambda = -\left(\frac{\partial \log \rho}{\partial \log T}\right)_P$$
$$x^3 + \frac{4}{9}\left(x^2 + x\right) = \frac{2}{9}\left(\frac{\lambda P}{\rho B^2}\right) \left(\nabla_{rad} - \nabla_{ad}\right), \quad H = min\left(\frac{P}{rhog}, \left(\frac{rP}{2\rho g}\right)^{1/2}\right), \quad v_{con} = \frac{1}{2}\,\alpha B\,x$$

### 3. Equation state and opacity

OPAL GN93 + Alexander opacities, OPAL 2001 state tables, converted for given Z to tables for log P, log U, log Cp,  $\lambda$ ,  $\nabla_{ad}$ ,  $\Gamma_1$ , log  $\kappa$  vs  $\left[\log T, \log(\rho/T^3)\right]$ . Interpolation by local cubics with continuous 1st derivatives. Minor species composition as in tables.

## 4. Surface layers - Atmosphere

Eddington grey atmosphere incorporated in model by imposing surface condition at R=r(N) where optical depth  $\tau=\tau_s\sim 0.001$ 

$$T^4(N) = \frac{L(N)}{4\pi\sigma R^2} \left(\tau_s + \frac{2}{3}\right), \qquad P(N) = \frac{GM(N)}{R^2} \frac{\tau_s}{\kappa(N)}$$

Photosphere determined by iterative interpolation to find the value of  $R_{ph}$  where  $T = T_{eff}$  with  $T_{eff}^4 = L(N)/(4\pi\sigma R_{ph}^2)$ .

## 5. Convective Core

Nearest mesh points relocated to core boundary  $(\nabla_{rad} = \nabla_{ad})$  and to overshoot boundary during successive iterations for structure and chemical evolution.

Chemical profiles outside *shrinking* cores smoothed linearly in  $M_r$  from  $M_c(t)$  to  $M_c(t+dt)$ Chemical overshooting only, extends mixed region by  $\beta \min(H, r_c)$ .

#### 6. Condensed Nuclear reaction network used in this code

 $R_{ij}X_iX_j$  = number of reactions/gm/sec of species *i* with species *j* Reaction rates  $R_{ij}$  and energy release  $Q_{ij}$  from NACRE with  $\nu, \beta$  decay from Bahcall. New analytic fit to weak-intermediate-strong screening.

$$\begin{split} &R_{11}:{}^{1}H\left(p,\nu\;e^{+}\right){}^{2}H(p,\;\gamma)\;{}^{3}He\\ &R_{33}:{}^{3}He\left({}^{3}He,\alpha\;2p\right){}^{4}He\\ &R_{43}:{}^{3}He\left(\alpha,\gamma\right){}^{7}Be\left(e^{-},\nu\right){}^{7}Li\left(p,\alpha\right){}^{4}He\\ &R_{121}:{}^{12}C\left(p,\gamma\right){}^{13}N\left(\;,e^{+}\;\nu\right){}^{13}C\\ &R_{131}:{}^{13}C\left(p,\gamma\right){}^{14}N\\ &R_{141}:{}^{14}N\left(p,\gamma\right){}^{15}O\left(\;,e^{+}\;\nu\right){}^{15}N\\ &R_{151}:{}^{15}N\left(p,\gamma\;\alpha\right){}^{12}C\\ &R_{151a}:{}^{15}N\left(p,\gamma\right){}^{16}O\\ &R_{161}:{}^{16}O\left(p,\gamma\right){}^{17}F\left(\;,e^{+}\;\nu\right){}^{17}O\\ &R_{171}:{}^{17}O\left(p,\gamma\;\alpha\right){}^{14}N \end{split}$$

#### 7. Chemical evolution equations

Mixing in convective regions is modelled as a diffusion process with the diffusion coefficient  $\nu_c$  determined by the MLT model of convection  $\nu_c = v_{con} \ell/3$ 

$$\frac{\partial X_k}{\partial t} = m_i \sum N_{ijk} R_{ij} X_i X_j + \frac{1}{\rho r^2} \frac{\partial}{\partial r} \left( \rho \nu_c r^2 \frac{\partial X_k}{\partial r} \right)$$

where  $N_{ijk}$  is the number of particles of species k produced in reaction  $R_{ij}$ 

#### Solving the chemical evolution equations for $X_k$

At mesh point i the diffusion term is discretised in conservative form as

$$-\frac{dt}{\rho r^2}\frac{\partial}{\partial r}\left(\rho\nu_c r^2\frac{\partial X_k}{\partial r}\right)_i = A_p[X_k(i+1) - X_k(i)] - A_m[X_k(i) - X_k(i-1)]$$

The evolution equations form a set of tridagonal equations for each k of the form

$$A_p(i)X_k(i+1) + A_{0k}(i)X_k(i) + A_m(i)X_k(i-1) = S_k(i), \quad i = 0, N; \quad k = 1, 9$$

where  $A_{0k}$  and  $S_k$  depend on the values of  $V_j, X_j$  at time t + dt and  $V_{0j}, X_{0j}$  at time t, The equations are solved using a 1<sup>st</sup> order implicit algorithm (of which there are many varities!).

The equations are solved sequentially; that is for each k we solve the system for i = 0, Nusing a tridiagonal matrix solver, and the set whole set is repeatedly solved with the updated  $X_j(i)$  until the solution for the  $X_k(i)$  has converged here defined as

$$\sum_{i} \left[\delta X_{1}(i)\right]^{2} + 10^{6} \sum_{i} \left[\delta X_{3}(i)\right]^{2} + 10^{4} \sum_{k \neq 1,3} \sum_{i} \left[\delta X_{k}(i)\right]^{2} < acc \ (\sim 10^{-10})$$

where  $\delta X_k$  is the difference in values of  $X_k$  between succesive iterations.

#### 8. Solving structure equations give $X_k$

The variables  $V_1(i) = r, V_2(i) = L_r, V_3(i) = \rho, V_4(i) = T$ ; all other state variables are known in terms of these variables and the values of  $X_k(i)$  and Z

The time derivatives  $\partial Q/\partial t$  are taken as  $1^{st}$  order implicit in time, and the differential equations are discretised to  $2^{nd}$  order in space in the form:

$$E_{1}(i) = [M_{i+1} - M_{i}] - \frac{1}{2} \left[ \left( \frac{dM_{r}}{dr^{3}} \right)_{i} + \left( \frac{dM_{r}}{dr^{3}} \right)_{i+1} \right] \left[ r_{i+1}^{3} - r_{i}^{3} \right]$$

$$E_{2}(i) = [L_{i+1} - L_{i}] - \frac{1}{2} \left[ \left( \frac{dL_{r}}{dM_{r}} \right)_{i} + \left( \frac{dL_{r}}{dM_{r}} \right)_{i+1} \right] [M_{i+1} - M_{i}]$$

$$E_{3}(i) = [T_{i+1} - T_{i}] - \frac{1}{2} \left[ \left( \frac{dT}{dr^{2}} \right)_{i} + \left( \frac{dT}{dr^{2}} \right)_{i+1} \right] \left[ r_{i+1}^{2} - r_{i}^{2} \right]$$

$$E_{4}(i) = \log \left( \frac{T_{i+1}}{T_{i}} \right) - \frac{1}{2} \left[ \nabla_{i+1} - \nabla_{i} \right] \log \left( \frac{P_{i+1}}{P_{i}} \right)$$

The equations are satisfied when  $E_k(i) = 0, k = 1, 4; i = 1, N-1$  plus the central boundary conditions  $(r = 0, L_r = 0 \text{ at } M_r = 0, \text{ and the surface boundary conditions given by the atmosphere (section 4 above).}$ 

The  $E_k(i)$  depend on the variables at  $V_j(i)$ ,  $V_j(i+1)$ , j = 1, 4. We iterate to find the values of the  $V_j$  that give  $E_k(i) = 0$  using a Newton-Raphson technique.

At any given iteration  $E_k(i) \neq 0$ . We find the derivatives of the  $E_k(i)$  wrt  $V_j(i), V_j(i+1)$ , and solve the linearised equations for corrections  $\delta V_j(i)$ 

$$\frac{\partial E_k(i)}{\partial V_j(i)}\delta V_j(i) + \frac{\partial E_k(i)}{\partial V_j(i+1)}\delta V_j(i+1) = -E_k(i)$$

which can be written as

 $Akj(i)\,\delta V_j(i) = -E_k(i)$ 

where A is a block diagonal matrix, the blocks being 8 x 4. This system is readily solved by elimination of the first 2 columns in each block, diagonalisation of the 4 x 4 square section of the block, and back substitution. This gives corrections  $\delta V_j(i)$  to be added to the  $V_j(i)$  This process is repeated until the solution has converged.

In practice we use log V rather than V and the solution is deemed to be converged when all corrections  $\delta V/V < acc \ (\sim 1/N^2)$ .

#### Results

The results for the ESTA comparison models are given in the poster by Montiero et al. Here we point out that relocating a mesh point at the boundary of convective cores and of any overshoot region results in a relatively smooth variation of the Brunt-Väisäla frequency. This is illustrated in the figure which shows the results of applying this code to Case 1.5 of the ESTA comparison exercise. This is for a star of  $M = 2.0 M_{\odot}, X_0 = 0.72, Z = 0.02$  with an overshoot parameter  $\beta = 0.15$  evolved to a central hydrogen abundance of  $X_c = 0.01$ .



